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# PARTIAL INTEGRAL OPERATORS OF FREDHOLM TYPE ON KAPLANSKY–HILBERT MODULE OVER $L_0$

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Dedicated to the 80th anniversary of Professor Stefan Grigorievich Samko

Abstract. The article studies some characteristic properties of self-adjoint partially integral operators of Fredholm type in the Kaplansky-Hilbert module  $L_0\left[L_2\left(\Omega_1\right)\right]$  over  $L_0\left(\Omega_2\right)$ . Some mathematical tools from the theory of Kaplansky-Hilbert module are used. In the Kaplansky-Hilbert module  $L_0\left[L_2\left(\Omega_1\right)\right]$  over  $L_0\left(\Omega_2\right)$  we consider the partially integral operator of Fredholm type  $T_1\left(\Omega_1\text{ and }\Omega_2\text{ are closed bounded sets in }\mathbb{R}^{\nu_1}\text{ and }\mathbb{R}^{\nu_2},\ \nu_1,\nu_2\in\mathbb{N}$ , respectively). The existence of  $L_0\left(\Omega_2\right)$  nonzero eigenvalues for any self-adjoint partially integral operator  $T_1$  is proved; moreover, it is shown that  $T_1$  has finite and countable number of real  $L_0\left(\Omega_2\right)$ -eigenvalues. In the latter case, the sequence  $L_0\left(\Omega_2\right)$ -eigenvalues is order convergent to the zero function. It is also established that the operator  $T_1$  admits an expansion into a series of  $\nabla_1$ -one-dimensional operators.

Key words: partial integral operator, Kaplansky–Hilbert module,  $L_0$ -eigenvalue.

Mathematical Subject Classification (2010): 45A05, 47A10, 47G10, 45P05, 45B05, 45C05.

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#### 1. Introduction

Linear equations and operators involving partial integrals appear in elasticity theory, continuum mechanics, aerodynamics and in PDE theory [1]. Self-adjoint partial integral operators arise in the theory of Schrodinger operators [2, 3]. Spectral properties of a discrete Schrodinger operator H are closely related (see [3, 4]) to the partial integral operators which participate in the presentation of operator H.

Let  $\Omega_1$  and  $\Omega_2$  be closed bounded subsets in  $\mathbb{R}^{\nu_1}$  and  $\mathbb{R}^{\nu_2}$ , respectively. Partial integral operator (PIO) of Fredholm type in the space  $L_p(\Omega_1 \times \Omega_2)$ ,  $p \geqslant 1$ , is an operator of the form [1]

$$T = T_0 + T_1 + T_2 + K, (1)$$

where operators  $T_0$ ,  $T_1$ ,  $T_2$  and K are defined by the following formulas

$$T_0 f(x, y) = k_0(x, y) f(x, y),$$

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$$T_1 f(x,y) = \int_{\Omega_1} k_1(x,s,y) f(s,y) ds,$$

$$T_2 f(x,y) = \int_{\Omega_2} k_2(x,t,y) f(x,t) dt,$$

$$K f(x,y) = \int_{\Omega_1} \int_{\Omega_2} k(x,y;s,t) f(s,t) ds dt.$$
(2)

Here  $k_0$ ,  $k_1$ ,  $k_2$  and k are given measurable functions on  $\Omega_1 \times \Omega_2$ ,  $\Omega_1^2 \times \Omega_2$ ,  $\Omega_1 \times \Omega_2^2$  and  $(\Omega_1 \times \Omega_2)^2$ , respectively, and all integrals have to be understood in the Lebesgue sense, where  $ds = d\mu_1(s), dt = d\mu_2(t), \mu_k(\cdot)$  — the Lebesgue measure on the  $\sigma$ -algebra of subsets  $\Omega_k$ , k = 1, 2.

Furthermore, some simple solvability conditions for the equations Tf = q were investigated by several authors (see, for example, [1] and its references). Spectral properties of the given operator has been studied in [1, 4, 5].

Nevertheless, the description of the spectra of self-adjoint PIOs with  $L_2$  kernels remains an open question. Difficulty of this problem is connected with non-compactness of the operators  $T_1$  and  $T_2$ . The article studies some characteristic properties of self-adjoint partially integral operators of Fredholm type in the Kaplansky-Hilbert module  $L_0[L_2(\Omega_1)]$  over  $L_0(\Omega_2)$ . The mathematical tools from the Kaplansky-Hilbert module is used as presented in [6].

The paper is organized as follows. In Section 3 we prove the existence of an  $L_0$ -eigenvalue for the PIO  $T_1$ .

In Section 4 we study existence of the countable consequence of real  $L_0$ -eigenvalues for PIO  $T_1$ . In Section 5 it is given the decomposition of the PIO  $T_1$  in series of  $\nabla_1$ -one-dimensional operators. In Section 5 (in section 6) is given decomposition of the PIO  $T_1$  (the PIO  $T_2$ ) in series of  $\nabla_{1}$ - ( $\nabla_{2}$ -) one-dimensional operators.

#### 2. Kaplansky–Hilbert Module over $L_0$

Recall some notions and results from the theory of Kaplansky-Hilbert modules (see [6]). Let  $(\Omega_k, \Sigma_k, \mu_k)$  be a space with complete finite measure  $\mu_k$ ,  $L_0(\Omega_k)$ -algebra of equivalence classes of all complex measurable functions on  $(\Omega_k, \Sigma_k, \mu_k)$ , where k = 1, 2. We denote by  $L_0[L_2(\Omega_1)]$  the set of equivalence classes of all complex measurable functions f(x,y) on  $\Omega_1 \times$  $\Omega_2$ , which satisfies the condition: the integral

$$\varphi(y) = \int_{\Omega_1} |f(x,y)|^2 d\mu_1(x)$$

exists for almost all  $y \in \Omega_2$  and  $\varphi \in L_0(\Omega_2)$ .

We consider the map  $\langle \cdot, \cdot \rangle_1 : L_0[L_2(\Omega_1)] \times L_0[L_2(\Omega_1)] \to L_0(\Omega_2)$  by rule

$$\langle f, g \rangle_1 = \int_{\Omega_1} f(s, y) \overline{g(s, y)} \, d\mu_1(s).$$

It is clear, that the map  $\langle \cdot, \cdot \rangle_1$  satisfies the conditions of  $L_0(\Omega_2)$ -valued inner product.

For each  $f \in L_0[L_2(\Omega_1)]$  we define  $L_0$ -norm:

$$||f||_1(\omega) = \sqrt{\langle f, f \rangle_1(\omega)}.$$

Then  $L_0[L_2(\Omega_1)]$  is Banach–Kantorovich space over  $L_0(\Omega_2)$  [6, 7]. Consequently, the space  $L_0[L_2(\Omega_1)]$  is Kaplansky-Hilbert module over  $L_0(\Omega_2)$  with the inner product  $\langle \cdot, \cdot \rangle_1(\omega)$ .

If for the map  $A: L_0[L_2(\Omega_1)] \to L_0[L_2(\Omega_1)]$  the equality  $A(\alpha \cdot f + \beta \cdot g) = \alpha \cdot Af + \beta \cdot Ag$  is hold for all  $\alpha, \beta \in L_0(\Omega_2)$ ,  $f, g \in L_0[L_2(\Omega_1)]$ , then A is called  $L_0(\Omega_2)$ -linear operator.

If for the  $L_0(\Omega_2)$ -linear operator A there exists  $C = C(\omega) \in L_0(\Omega_2)$  such that,  $||Af||_1(\omega) \leq C(\omega)||f||_1(\omega)$  for all  $f \in L_0[L_2(\Omega_1)]$ , then A is called  $L_0(\Omega_2)$ -bounded operator.

For each  $L_0(\Omega_2)$ -linear  $L_0(\Omega_2)$ -bounded operator A we define  $L_0(\Omega_2)$ -norm by the rule

$$||A||_1 = ||A||_1(\omega) = \sup\{||Af||_1(\omega) : ||f||_1 \le \mathbf{e}\}.$$

We say the net  $(\xi_{\alpha})_{\alpha \in A} \subset L_0(\Omega_2)$  (o)-converges to the element  $\xi \in L_0(\Omega_2)$ , whenever there is a decreasing net  $(e_{\beta})_{\beta \in B} \subset L_0(\Omega_2)$  such that  $\inf\{e_{\beta} : \beta \in B\} = \theta$  and for each  $\beta \in B$  there is an index  $\alpha(\beta) \in A$  with  $|\xi_{\alpha} - \xi| \leq e_{\beta}$  for all  $\alpha \in A : \alpha(\beta) \leq \alpha$ . In this case, the element  $\xi$  is called (o)-limit of the set  $(\xi_{\alpha})_{\alpha \in A}$  and we write  $\xi = (o)$ -  $\lim \xi_{\alpha}$ .

We know [8], that the (o)-converges of the net  $(\xi_{\alpha})_{\alpha \in A} \subset L_0(\Omega_2)$  to the element  $\xi$  is equivalent to converges almost everywhere to the element  $\xi$  of the net  $(\xi_{\alpha})_{\alpha \in A} \subset L_0(\Omega_2)$ .

The net  $(f_{\alpha})_{\alpha \in A}$  in  $L_0[L_2(\Omega_1)]$  is called (bo)-converging to  $f \in L_0[L_2(\Omega_1)]$ , if (o)-  $\lim ||f_{\alpha} - f||_1 = \theta$  in  $L_0(\Omega_2)$ .

Let  $\Lambda_2$  be the Boolean algebra of idempotents in  $L_0(\Omega_2)$ . If  $(f_{\alpha})_{\alpha \in A} \subset L_0[L_2(\Omega_1)]$  and  $(\pi_{\alpha})_{\alpha \in A}$  is a partition of the unit in  $\Lambda_2$ , then the series  $\sum_{\alpha} \pi_{\alpha} \cdot f_{\alpha}$  (bo)-converges in  $L_0[L_2(\Omega_1)]$  and its sum is called the mixing of  $(f_{\alpha})_{\alpha \in A}$  with respect to  $(\pi_{\alpha})_{\alpha \in A}$ . We denote this sum by  $\min(\pi_{\alpha}f_{\alpha})$ . A subset  $K \subset L_0[L_2(\Omega_1)]$  is called cyclic, if  $\min(\pi_{\alpha}f_{\alpha}) \in K$  for each  $(f_{\alpha})_{\alpha \in A} \subset K$  and any partition of the unit  $(\pi_{\alpha})_{\alpha \in A}$  in  $\Lambda_2$ . A subset  $K \subset L_0[L_2(\Omega_1)]$  is called cyclically compact, if K is cyclic and every net in K has a cyclic subset that (bo)-converges to some point of K. A subset is called relatively cyclically compact, if it is contained in a cyclically compact set.

A  $L_0$ -linear operator in  $L_0[L_2(\Omega_1)]$  is called *cyclically compact*, if for every  $L_0$ -bounded set B in  $L_0[L_2(\Omega_1)]$  the set A(B) is relatively cyclically compact in  $L_0[L_2(\Omega_1)]$ .

Let  $T_1$  be an operator in the Kaplansky–Hilbert module  $L_0[L_2(\Omega_1)]$  over  $L_0(\Omega_2)$  given by the formula

$$(T_1 f)(x, y) = \int_{\Omega_1} k_1(x, s, y) f(s, y) d\mu_1(s).$$
(3)

Here,  $k_1(x, s, y)$  is a measurable function on  $\Omega_1^2 \times \Omega_2$ .

Let the kernel  $k_1(x, s, y)$  of the integral operator  $T_1$  satisfy the condition

$$\int_{\Omega_1} \int_{\Omega_1} |k_1(x, s, y)|^2 d\mu_1(s) d\mu_1(x) \in L_0(\Omega_2).$$
(4)

Then, the operator  $T_1$  with values in  $L_0(\Omega_2)$  is linear and bounded on  $L_0[L_2(\Omega_1)]$ .

Also, let the kernel  $k_1(x, s, y)$  satisfy the condition:

$$k_1(x, s, y) = \overline{k_1(s, x, y)}.$$

Then the operator  $T_1$  is a self-adjoint operator on the Kaplansky-Hilbert module  $L_0[L_2(\Omega_1)]$ , i. e.,

$$\langle T_1 f, g \rangle_1 = \langle f, T_1 g \rangle_1.$$

A system  $\{f_{\alpha}(x,y)\}\subset L_0[L_2(\Omega_1)]$  is  $\nabla_1$ -orthogonal system, if  $\langle f_{\alpha}, f_{\beta}\rangle_1=\theta, \ \alpha\neq\beta$ . A  $\nabla_1$ -orthogonal system  $\{f_{\alpha}(x,y)\}\subset L_0[L_2(\Omega_1)]$  is said to be  $\nabla_1$ -orthonormal system, if  $\langle f_{\alpha}, f_{\alpha}\rangle_1=\mathbf{e}$ .

Note that, the PIO  $T_1$  is a good example for cyclically compact operators on Kaplansky–Hilbert module [7].

#### 3. $L_0$ -Eigenvalue of the Partial Integral Operator $T_1$

In this section we prove the existence of an  $L_0$ -eigenvalue for the PIO. Put  $\mathscr{H} = L_0[L_2(\Omega_1)]$ .

**Theorem 3.1.** The partial integral operator  $T_1$  has non zero  $L_0$ -eigenvalue.

< Put

$$\mathscr{D}_0 = \left\{ \omega \in \Omega_2 : \int_{\Omega_1} \int_{\Omega_1} |k_1(x, s, \omega)|^2 d\mu_1(x) d\mu_1(s) > 0 \right\}.$$

Then  $\mu_2(\mathcal{D}_0) = 0$ . For each  $f \in \mathcal{H}$ ,  $f \neq \theta$  we define subset  $\operatorname{supp}_{\Omega_2}(f)$  with positive measure by the following equality

$$\sup_{\Omega_2}(f) = \{ \omega \in \Omega_2 : \langle f, f \rangle_1(\omega) \neq 0 \}.$$

Let  $f_0 \in \mathcal{H}$ ,  $||f_0||_1(\omega) \neq 0$  for all  $\omega \in \mathcal{D}_0$  and  $T_1 f_0 \neq \theta$ . It is clear, that  $T_1^n f_0 \neq \theta$  for all  $n \in \mathbb{N}$ , as: if

$$T_1^k f_0 \neq \theta$$
,  $T_1^{k+1} f_0 = \theta$ , for some  $k \geqslant 1$ 

then we get a contradiction

$$\theta = \left\langle T_1^{k+1} f_0, T_1^{k-1} f_0 \right\rangle_1(\omega) = \left\langle T_1^k f_0, T_1^k f_0 \right\rangle_1(\omega) \neq \theta.$$

We construct two sequences  $\{\widetilde{f}_k(x,\omega)\}_{k\in\mathbb{N}_0}$ ,  $\{f_k(x,\omega)\}_{k\in\mathbb{N}_0}$  of functions from the Kaplansky–Hilbert module  $L_0[L_2(\Omega_1)]$  ( $\mathbb{N}_0=\mathbb{N}\cup 0$ ):

$$\widetilde{f}_k(x,\omega) = \begin{cases} \frac{f_k(x,\omega)}{\|f_k\|_1(\omega)}, & x \in \Omega_1, \ \omega \in \sup_{\Omega_2}(f_k), \\ 0, & x \in \Omega_1, \ \omega \in \Omega_2 \setminus \sup_{\Omega_2}(f_k), \end{cases}$$

$$f_{k+1}(x,\omega) = (T_1 \widetilde{f}_k)(x,\omega).$$

It follows from [9] that

$$||f_k||_1(\omega) \leqslant ||f_{k+1}||_1(\omega), \quad k \in \mathbb{N}, \tag{5}$$

and

$$||f_{k+1}||_1(\omega) \cdot ||f_k||_1(\omega) = \langle f_{k-1}, f_{k+1} \rangle_1(\omega) = \langle f_{k+1}, f_{k-1} \rangle_1(\omega), \quad k \in \mathbb{N}.$$
 (6)

On the other hand

$$||T_1\widetilde{f}_{k-1}||_1(\omega) \leqslant ||T_1||_1(\omega), \quad k \in \mathbb{N},$$

where  $||T_1||_1(\omega) \in L_0(\Omega_2)$  is the  $L_0(\Omega_2)$  valued norm of the PIO  $T_1$ . Consequently,

$$||f_k||_1(\omega) \leqslant ||T_1||_1(\omega), \quad k \in \mathbb{N}.$$

Thus, for almost all  $\omega \in \Omega_2$  the sequence  $\{\|f_k\|_1(\omega)\}_{k\in\mathbb{N}}$  has a finite limit  $\lambda(\omega) \geq 0$ , i. e.,

$$\lim_{k \to \infty} ||f_k||_1(\omega) = \lambda(\omega),\tag{7}$$

for almost all  $\omega \in \Omega_2$ . We have  $\lambda(\omega) \in L_0(\Omega_2)$ , as  $||f_k||_1(\omega) \in L_0(\Omega_2)$ ,  $k \in \mathbb{N}$ . From the relation (5) it follows that  $\lambda \neq \theta$ . Now, we define the family of integral operators  $\{T_1(\omega)\}$  on  $L_2(\Omega_1)$  by

$$T_1(\omega)\varphi(x) = \int_{\Omega_1} k_1(x, s, \omega)\varphi(s) d\mu_1(s), \quad \varphi \in L_2(\Omega_1), \ \omega \in \Omega_2.$$

Then,  $T_1(\omega)$  is a compact operator on  $L_2(\Omega_1)$  for almost all  $\omega \in \Omega_2$ . By the compactness of the operator  $T_1(\omega)$  there exists subsequence  $\widetilde{f}_{n_i}(x,\omega)$  such that  $f_{n_i+1}(x,\omega) = T_1(\omega)\widetilde{f}_{n_i}(x,\omega)$  has a limit  $g(x,\omega)$  in the  $L_0$ -norm  $\|\cdot\|_1$ . It is clear  $g \in \mathscr{H}$  and  $g \neq \theta$ . Analogously, for each sequence

$$f_{n_i+2}(x,\omega) = T_1(\omega)\widetilde{f}_{n_i+1}(x,\omega), \quad f_{n_i+3}(x,\omega) = T_1(\omega)\widetilde{f}_{n_i+2}(x,\omega)$$

we obtain  $f_{n_i+2} \to h \in \mathscr{H}$  and  $f_{n_i+3} \to \widetilde{h} \in \mathscr{H}$  by the  $L_0$ -norm  $\|\cdot\|_1$ . Using the relations (6), (7) we obtain

$$\begin{split} \|\widetilde{h} - g\|_1^2(\omega) &= \lim_{k \to \infty} \|f_{n_k+3} - f_{n_k+1}\|_1^2(\omega) \\ &= \lim_{k \to \infty} \left\{ \|f_{n_k+3}\|_1^2(\omega) + \|f_{n_k+1}\|_1^2(\omega) - \langle f_{n_k+3}, f_{n_k+1} \rangle_1(\omega) - \langle f_{n_k+1}, f_{n_k+3} \rangle_1(\omega) \right\} = 0 \end{split}$$

for almost all  $\omega \in \Omega_2$  and so  $\tilde{h} = g$ . On the other hand, from the equalities

$$f_{n_k+2}(x,\omega) = \begin{cases} \frac{(T_1(\omega)f_{n_k+1})(x,\omega)}{\|f_{n_k+1}\|_1(\omega)}, & x \in \Omega_1, \ \omega \in \sup_{\Omega_2} (f_{n_k+1}), \\ 0, & x \in \Omega_1, \ \omega \in \Omega_2 \setminus \sup_{\Omega_2} (f_{n_k+1}), \end{cases}$$

$$f_{n_k+3}(x,\omega) = \begin{cases} \frac{(T_1(\omega)f_{n_k+2})(x,\omega)}{\|f_{n_k+2}\|_1(\omega)}, & x \in \Omega_1, \ \omega \in \sup_{\Omega_2} (f_{n_k+2}), \\ 0, & x \in \Omega_1, \ \omega \in \Omega_2 \setminus \sup_{\Omega_2} (f_{n_k+2}) \end{cases}$$

we have

$$||f_{n_k+1}||_1(\omega) \cdot f_{n_k+2}(x,\omega) = (T_1(\omega)f_{n_k+1})(x,\omega), \quad \omega \in \Omega_2,$$
(8)

$$||f_{n_k+2}||_1(\omega) \cdot f_{n_k+3}(x,\omega) = (T_1(\omega)f_{n_k+2})(x,\omega), \quad \omega \in \Omega_2.$$
 (9)

It is clear that

$$\lim_{k \to \infty} \|f_{n_k+1}\|_1(\omega) = \|g\|_1(\omega) = \lim_{k \to \infty} \|f_{n_k+2}\|_1(\omega) = \|h\|_1(\omega)$$
$$= \lim_{k \to \infty} \|f_{n_k+3}\|_1(\omega) = \|\widetilde{h}\|_1(\omega) = \lambda(\omega).$$

From the equalities (8), (9) it follows that

$$\lambda(\omega) \cdot h(x,\omega) = T_1(\omega)g(x,\omega), \quad \lambda(\omega) \cdot \widetilde{h}(x,\omega) = T_1(\omega)h(x,\omega),$$

i. e.,

$$(T_1g)(x,y) = \lambda(y) \cdot h(x,y), \quad (T_1h)(x,y) = \lambda(y) \cdot g(x,y).$$

Hence it follows that

$$T_1(h+q)(x,y) = \lambda(y) \cdot (h+q)(x,y), \quad T_1(h-q)(x,y) = -\lambda(y) \cdot (h-q)(x,y).$$

We know, that  $h \neq \theta$ ,  $g \neq \theta$ . Hence we can conclude that:  $h + g \neq \theta$  or  $h - g \neq \theta$ . It means that the function  $\lambda(y)$  is an  $L_0$ -eigenvalue of the PIO  $T_1$ .  $\triangleright$ 

# 4. Spectral Properties of the Partial Integral Operator $T_1$ on the Kaplansky-Hilbert Module $L_0[L_2(\Omega_1)]$

**Theorem 4.1.** For a PIO  $T_1$  the following function  $\lambda_0(\omega) = \sup_{\|g\|_1 = \mathbf{e}} |\langle T_1 g, g \rangle_1(\omega)|$  is nonzero and either  $+\lambda_0(\omega)$  or  $-\lambda_0(\omega)$  is  $L_0$ -eigenvalue of the  $T_1$ .

⊲ Put

$$\Omega_0 = \left\{ \omega \in \Omega_2 : \int_{\Omega_1} \int_{\Omega_1} |k_1(x, s, \omega)|^2 d\mu_1(x) d\mu_1(s) > 0 \right\}.$$

From the  $T_1 \neq \theta$  it follows that  $\lambda_0(\omega) \neq 0$  for all  $\omega \in \Omega_0$ , i. e.,  $\lambda_0 \neq \theta$ . It is clear, that there is a sequence of  $\nabla_1$ -normal functions  $\{g_n\}_{n=1}^{\infty}$ , in which a limit exists

(o)-
$$\lim_{n\to\infty} \langle T_1 g_n, g_n \rangle_1(\omega) = \lambda(\omega),$$

and  $\lambda(\omega)$  is a real function on  $\Omega_2$ , where  $\lambda(\omega) = +\lambda_0(\omega)$  or  $-\lambda_0(\omega)$ . Consequently,  $\lambda_0 \in L_0(\Omega_2)$  and  $\operatorname{supp}(\lambda) = \Omega_0$ .

By cyclical compactness of the PIO  $T_1$  there exists a subsequence  $\{g_{n_i}\}_{i=1}^{\infty}$  with

(bo)-
$$\lim_{k \to \infty} (T_1 g_{n_k})(x, y) = h(x, y).$$
 (10)

Clearly,  $\operatorname{supp}_{\Omega_2}(h) = \Omega_0$ . From the equality

$$||T_1 g_{n_k} - \lambda \cdot g_{n_k}||_1^2 = ||T_1 g_{n_k}||_1^2 - 2\lambda \cdot \langle T_1 g_{n_k}, g_{n_k} \rangle_1 + \lambda^2$$

we obtain

(o)-
$$\lim_{k \to \infty} ||T_1 g_{n_k} - \lambda \cdot g_{n_k}||_1^2 = ||h||_1^2 - \lambda^2.$$
 (11)

However,

$$||T_1g_{n_k}||_1(\omega) \leqslant \lambda_0(\omega) \cdot ||g_{n_k}||_1(\omega) = |\lambda(\omega)|.$$

Therefore,

$$||h||_1(\omega) \leqslant |\lambda(\omega)|.$$

From this and (11) we have  $||h||_1(\omega) = |\lambda(\omega)|$ . Thus,

(o)- 
$$\lim_{k \to \infty} ||T_1 g_{n_k} - \lambda \cdot g_{n_k}||_1 = \theta.$$
 (12)

Hence, it follows that

$$T_1 f_0 = \lambda \cdot f_0$$

where

$$f_0(x,\omega) = \begin{cases} \frac{h(x,\omega)}{\lambda(\omega)}, & x \in \Omega_1, \ \omega \in \text{supp}(\lambda), \\ 0, & x \in \Omega_1, \ \omega \in \Omega_2 \setminus \text{supp}(\lambda). \end{cases}$$

Put

$$\pi_0(\omega) = \begin{cases} 1, & \omega \in \text{supp}(\lambda), \\ 0, & \Omega_2 \\ 0, & \omega \in \Omega_2 \setminus \text{supp}(\lambda). \end{cases}$$

Remark 4.1. Every element  $\zeta \in L_0(\Omega_2)$ ,  $\pi_0 \zeta = \lambda$  is  $L_0$ -eigenvalue of the PIO  $T_1$ .

**Theorem 4.2.** The PIO  $T_1$  has a finite or countable sequence of  $\nabla_1$ -orthonormal eigenfunctions

$$\phi_1(x,y), \phi_2(x,y), \ldots, \phi_n(x,y), \ldots$$

corresponding to a system of real nonzero  $L_0$ -eigenvalues

$$\lambda_1(\omega), \lambda_2(\omega), \ldots, \lambda_n(\omega), \ldots,$$

where

$$|\lambda_1(\omega)| \geqslant |\lambda_2(\omega)| \geqslant \ldots \geqslant |\lambda_n(\omega)| \geqslant \ldots$$

Moreover, for each  $f(x,y) \in L_0[L_2(\Omega_1)]$  the equality

$$||f||_1^2(\omega) = (o) - \sum_{k=1}^{\infty} |\langle f, \phi_k \rangle_1(\omega)|^2$$

holds.

 $\lhd$  Put  $\mathscr{H}_1 = \mathscr{H}$  and  $T_1^{(1)} = T_1$ . By the Theorem 4.1 there is such element  $\phi_1(x,y) \in \mathscr{H}_1$  that  $T_1^{(1)}\phi_1 = \lambda_1 \cdot \phi_1$ , where  $\lambda_1$  is a real function on  $\Omega_2$  and  $\lambda_1(\omega) = \pm \sup_{\|g\|_1 = \mathbf{e}} |\langle T_1 g, g \rangle_1(\omega)|$ . We define the Kaplansky-Hilbert submodule  $\mathscr{H}_2 = \mathscr{H}_1 \ominus_1 \{\phi_1\}$ . It is clear that if  $f \in \mathscr{H}_2$ , then  $T_1^{(1)}f \in \mathscr{H}_2$  from the equality  $\langle f, \phi_1 \rangle_1 = \theta$  it follows that

$$\langle T_1^{(1)} f, \phi_1 \rangle_1 = \langle f, T_1 \phi_1 \rangle_1 = \langle f, \lambda_1 \cdot \phi_1 \rangle_1 = \theta.$$

We define an operator  $T_1^{(2)}$  on the  $\mathcal{H}_2$  by

$$T_1^{(2)}f = T_1^{(1)}f, \quad f \in \mathcal{H}_2.$$

The operator  $T_1^{(2)}$  is a selfadjoint PIO on the  $\mathscr{H}_2$ . If  $T_1^{(2)} \neq \theta$ , then we apply Theorem 4.1 to the operator  $T_1^{(2)}$  and find an element  $\phi_2(x,y) \in \mathscr{H}_2$  such that  $T_1^{(2)}\phi_2 = \lambda_2 \cdot \phi_2$ , where  $\lambda_2$  is a real function on  $\Omega_2$  and  $\lambda_2(\omega) = \pm \sup_{g \in \mathscr{H}_2, \|g\|_1 = \mathbf{e}} \left| \left\langle T_1^{(2)}g, g \right\rangle_1(\omega) \right|$ . As  $\phi_2(x,y) \in \mathscr{H}_2$ ,  $\|\phi_2\|_1 = \mathbf{e}$ , we have  $\langle \phi_2, \phi_1 \rangle_1 = \theta$ . Therefore,

$$|\lambda_2(\omega)| = \sup_{g \in \mathcal{H}_2, \|g\|_1 = \mathbf{e}} \left| \left\langle T_1^{(1)} g, g \right\rangle_1(\omega) \right| \leqslant \sup_{g \in \mathcal{H}_1, \|g\|_1 = \mathbf{e}} \left| \left\langle T_1^{(1)} g, g \right\rangle_1(\omega) \right| = |\lambda_1(\omega)|.$$

Continuing this process we obtain a sequence of Kaplansky–Hilbert submodules  $\mathscr{H}_{k+1} = \mathscr{H}_k \ominus_1 \{\phi_k\}$ , where  $\phi_k \in \mathscr{H}_k$  are eigenfunctions of the PIO  $T_1$  with  $T_1\phi_k = \lambda_k \cdot \phi_k$ .

If  $T_1^{(n)}$  is a zero operator for some  $n \in \mathbb{N}$  then we obtain the finite system  $\nabla_1$ -orthonormal eigenfunctions  $\phi_1(x,y), \phi_2(x,y), \dots, \phi_{n-1}(x,y)$  corresponding to the system of nonzero  $L_0$ -eigenvalues  $\lambda_1(\omega), \lambda_2(\omega), \dots, \lambda_{n-1}(\omega)$ , such that

$$|\lambda_1(\omega)| \geqslant |\lambda_2(\omega)| \geqslant \ldots \geqslant |\lambda_{n-1}(\omega)|$$

and

$$|\lambda_k(\omega)| = \sup_{g \in \mathscr{H}_k, \|g\|_1 = \mathbf{e}} |\langle T_1^{(1)} g, g \rangle_1(\omega)|.$$

If  $T_1^{(n)} \neq \theta$  for each  $n \in \mathbb{N}$  then we obtain an infinite system  $\nabla_1$ -orthonormal eigenfunctions  $\{\phi_k\}_{k=1}^{\infty}$  corresponding to the system of  $L_0$ -eigenvalues  $\lambda_k \neq \theta$ . However, the equality

$$T_1(\omega)\phi_k(x,\omega) = \lambda_k(\omega) \cdot \phi_k(x,\omega), \quad k \in \mathbb{N},$$

is correct for almost all  $\omega \in \Omega_2$ . It follows that  $\lim_{k\to\infty} \lambda_k(\omega) = 0$  for almost all  $\omega \in \Omega_2$ , because  $T_1(\omega)$  is a compact operator for almost all  $\omega \in \Omega_2$ .

Let  $f = T_1 h$ ,  $h \in \mathcal{H}$  and  $g = h - \sum_{k=1}^m \langle h, \phi_k \rangle_1 \cdot \phi_k$ . Here m is the number of eigenfunctions of the system  $\{\phi_k\}$  when the system  $\{\phi_k\}$  is a finite set, and m is equal to arbitrary natural number otherwise. By the equality

$$\langle g, \phi_k \rangle_1 = \theta, \quad k \in \{1, 2, \dots, m\}$$

we have  $g \in \mathcal{H}_{m+1}$ . Consequently, we have

$$||T_1g||_1^2(\omega) \le ||T_1^{(m+1)}||^2(\omega) \cdot ||g||_1^2(\omega),$$

i. e.,

$$\left\| T_1 h - \sum_{k=1}^{m} \langle h, \phi_k \rangle_1 \cdot T_1 \phi_k \right\|_1^2 (\omega) \leqslant \left\| T_1^{(m+1)} \right\|^2 (\omega) \cdot \|g\|_1^2 (\omega). \tag{13}$$

We have  $\langle h, \phi_k \rangle_1 \cdot T_1 \phi_k = \langle T_1 h, \phi_k \rangle_1 \cdot \phi_k$  and  $\|g\|_1 \leqslant \|h\|_1$ . Hence by the inequality (13) we obtain

$$\left\| f - \sum_{k=1}^{m} \langle f, \phi_k \rangle_1 \cdot \phi_k \right\|_1^2(\omega) \leqslant \left\| T_1^{(m+1)} \right\|_1^2(\omega) \cdot \|h\|_1^2(\omega). \tag{14}$$

If the number of elements of the system  $\{\phi_k\}$  is equal to m then  $T_1^{(m+1)}=\theta$  and we have

$$f = \sum_{k=1}^{m} \langle f, \phi_k \rangle_1 \cdot \phi_k.$$

If the sequence  $\{\phi_k\}$  is infinite then from the inequality (14) it follows that

$$\left\| f - \sum_{k=1}^{m} \langle f, \phi_k \rangle_1 \cdot \phi_k \right\|_1^2(\omega) \leqslant \lambda_{m+1}^2(\omega) \cdot \|h\|_1^2(\omega),$$

i. e.,

$$\theta \leqslant \|f\|_1^2(\omega) - \sum_{k=1}^m |\langle f, \phi_k \rangle_1(\omega)|^2 \leqslant \lambda_{m+1}^2(\omega) \cdot \|h\|_1^2(\omega).$$

Thus as  $m \to \infty$ , we get

$$||f||_1^2(\omega) = (o) - \sum_{k=1}^{\infty} |\langle f, \phi_k \rangle_1(\omega)|^2. >$$

### 5. Decomposition of the Partial Integral Operator $T_1$ in Series of $\nabla_1$ -One-Dimensional Operators

DEFINITION 5.1. If for an operator  $A: \mathcal{H} \to \mathcal{H}$  there are  $\nabla_1$ -orthonormal functions  $\{\phi_k\}_{k=1}^n \subset \mathcal{H}$  and some system of functions  $\{g_k\}_{k=1}^n \subset \mathcal{H}$ , such that

$$Af = \sum_{k=1}^{n} \langle f.g_k \rangle_1 \cdot \phi_k, \quad f \in \mathcal{H}$$

then the operator A is called the  $\nabla_1$ -n-dimensional operator, here  $\mathcal{H} = L_0[L_2(\Omega_1)]$ .

**Theorem 5.1.** For the PIO  $T_1$  there is a system of  $\nabla_1$ -orthonormal functions  $\{\phi_k(x,y)\}$  and a sequence of real  $L_0(\Omega_2)$ -eigenvalues  $\lambda_k(\omega)$  such that for all  $h \in L_0[L_2(\Omega_1)]$  the following conditions hold:

1°.  $h = h_0 + (bo) - \sum_{k=1}^{\infty} \langle h, \phi_k \rangle_1 \cdot \phi_k, \ h_0 \in Ker(T_1).$ 

2°.  $T_1h = (bo) - \sum_{k=1}^{\infty} \lambda_k \cdot \langle h, \phi_k \rangle_1 \cdot \phi_k$ .

 $3^{\circ}$ .  $|\lambda_k(\omega)| \geqslant |\lambda_{k+1}(\omega)|, k \in \mathbb{N}$ .

 $4^{\circ}$ . (o)- $\lim_{k\to\infty} \lambda_k = \theta$ .

 $\triangleleft$  By Theorem 4.2 there are a system of  $\nabla_1$ -orthonormal functions  $\{\phi_k(x,y)\}$  and a sequence of  $L_0$ -eigenvalues  $\lambda_k(\omega)$  such that  $T\phi_k = \lambda_k \cdot \phi_k$  and for each  $f = T_1 h$  we get the equality

$$f = (bo)$$
-  $\sum_{k=1}^{\infty} \langle f, \phi_k \rangle_1 \cdot \phi_k$ ,

where  $\langle f, \phi_k \rangle_1 = \lambda_k \cdot \langle h, \phi_k \rangle_1$ .

Thus, for all  $h \in \mathcal{H}$ 

$$T_1 h = (bo) - \sum_{k=1}^{\infty} \lambda_k \cdot \langle h, \phi_k \rangle_1 \cdot \phi_k.$$

If we denote  $h_0 = h - (bo) - \sum_{k=1}^{\infty} \langle h, \phi_k \rangle_1 \cdot \phi_k$ , then

$$h = h_0 + (bo) - \sum_{k=1}^{\infty} \langle h, \phi_k \rangle_1 \cdot \phi_k, \quad T_1 h_0 = \theta.$$

The properties  $3^{\circ}$  and  $4^{\circ}$  follows from the Theorem 4.2. Theorem 5.1 can also be proven by using Theorem 3.5 in the article of A. G. Kusraev [10].  $\triangleright$ 

**Theorem 5.2.** For all positive functions  $\varepsilon(\omega) \in L_0(\Omega_2)$ ,  $\mu_2(\Omega_2 \setminus \text{supp}(\varepsilon)) = 0$  there exist a  $\nabla_1$ -finite dimensional operator  $\mathscr{T}_1^{\varepsilon}$  on the Kaplansky–Hilbert module  $L_0[L_2(\Omega_1)]$ , such that  $||T_1 - \mathscr{T}_1^{\varepsilon}||_1(\omega) < \varepsilon(\omega)$ .

 $\triangleleft$  By the Theorem 5.1 there is a system of  $\nabla_1$ -orthonormal functions  $\{\phi_k(x,y)\}$  and a sequence of  $L_0$ -eigenvalues  $\lambda_k(\omega)$  for which the properties 1°-4° hold. We define the  $\nabla_1$ -finite dimensional operator  $\mathscr{T}_1^{\varepsilon}$ :

$$\mathscr{T}_1^{\varepsilon} h = \sum_{k=1}^n \lambda_k \cdot \langle h, \phi_k \rangle_1 \cdot \phi_k.$$

It follows that

$$||T_1h - \mathcal{J}_1^{\varepsilon}h||_1^2(\omega) \leqslant \lambda_{n+1}^2(\omega) \cdot \left\{ |\langle h, \phi_{k+1} \rangle_1(\omega)|^2 + |\langle h, \phi_{k+2} \rangle_1(\omega)|^2 + \ldots \right\} \leqslant \lambda_k^2(\omega) ||h||_1^2(\omega).$$

Hence, for  $|\lambda_{n+1}(\omega)| < \varepsilon(\omega)$  we have  $||T_1 - \mathcal{I}_1^{\varepsilon}||_1(\omega) < \varepsilon(\omega)$ .  $\triangleright$ 

# 6. Decomposition of the Partial Integral Operator $T_2$ in Series of $\nabla_2$ -One-Dimensional Operators

We denote by  $L_0[L_2(\Omega_2)]$  the set of equivalence classes of all complex measurable functions f(x,y) on  $\Omega_1 \times \Omega_2$ , which satisfied the condition: the integral

$$\psi(x) = \int_{\Omega_2} |f(x,y)|^2 d\mu_2(y)$$

exist for almost all  $x \in \Omega_1$  and  $\psi \in L_0(\Omega_1)$ .

We define  $L_0(\Omega_1)$ -valued inner product on  $L_0[L_2(\Omega_2)]$  by

$$\langle f, g \rangle_2 = \int_{\Omega_1} f(x, t) \overline{g(x, t)} \, d\mu_2(t).$$

For each  $f \in L_0[L_2(\Omega_2)]$  we define  $L_0$ -norm:  $||f||_2(v) = \sqrt{\langle f, f \rangle_2(v)}$ . Then  $L_0[L_2(\Omega_2)]$  is a Banach-Kantorovich space over  $L_0(\Omega_1)$ . Consequently, the space  $L_0[L_2(\Omega_2)]$  is a Kaplansky-Hilbert module over  $L_0(\Omega_1)$  with the inner product  $\langle \cdot, \cdot \rangle_2(v)$ .

Let  $T_2$  be an operator in the Kaplansky-Hilbert module  $L_0[L_2(\Omega_2)]$  over  $L_0(\Omega_1)$  given by the formula

$$(T_2 f)(x, y) = \int_{\Omega_2} k_2(x, t, y) f(x, t) d\mu_2(t).$$
 (15)

Here,  $k_2(x, t, y)$  is measurable function on  $\Omega_1 \times \Omega_2^2$ .

Assume that the kernel  $k_2(x, s, y)$  of the integral operator  $T_2$  satisfies the condition

$$\int_{\Omega_2} \int_{\Omega_2} |k_2(x,t,y)|^2 d\mu_2(t) d\mu_2(y) \in L_0(\Omega_1).$$

Then, the operator  $T_2$  is linear and  $L_0(\Omega_1)$ -bounded operator on  $L_0[L_2(\Omega_2)]$ . If the kernel  $k_2(x,s,y)$  satisfy of the condition  $k_2(x,t,y)=k_2(x,y,t)$ , then the operator  $T_2$  is a self-adjoint operator on the Kaplansky-Hilbert module  $L_0[L_2(\Omega_2)]$ , i. e.,

$$\langle T_2 f, g \rangle_2 = \langle f, T_2 g \rangle_2.$$

A system  $\{f_{\alpha}(x,y)\}\in L_0[L_2(\Omega_2)]$  is said  $\nabla_2$ -orthogonal system, if  $\langle f_{\alpha},f_{\beta}\rangle_2=\theta,\,\alpha\neq\beta$ . A  $\nabla_2$ orthogonal system  $\{f_{\alpha}(x,y)\}\subset L_0[L_2(\Omega_2)]$  is said  $\nabla_2$ -orthonormal system, if  $\langle f_{\alpha},f_{\alpha}\rangle_2=\mathbf{e}$ .

Note that, the PIO  $T_2$  is cyclically compact on the Kaplansky-Hilbert module  $L_0[L_2(\Omega_2)]$  [7].

**Theorem 6.1.** For the PIO  $T_2$  there is a system of  $\nabla_2$ -orthonormal functions  $\{\psi_k(x,y)\}$ and a sequence of real  $L_0(\Omega_1)$ -eigenvalues  $\zeta_k(v)$  such that, for all  $h \in L_0[L_2(\Omega_2)]$  the following

- 1°.  $h = h_0 + (bo) \sum_{k=1}^{\infty} \langle h, \psi_k \rangle_1 \cdot \psi_k, h_0 \in Ker(T_2);$ 2°.  $T_2 h = (bo) \sum_{k=1}^{\infty} \zeta_k \cdot \langle h, \psi_k \rangle_1 \cdot \psi_k, \text{ where}$
- $3^{\circ}$ .  $|\zeta_k(v)| \geqslant |\zeta_{k+1}(v)|, k \in \mathbb{N};$
- $4^{\circ}$ . (o)- $\lim_{k\to\infty}\zeta_k=\theta$ .

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### ЧАСТИЧНО ИНТЕГРАЛЬНЫЕ ОПЕРАТОРЫ ТИПА ФРЕДГОЛЬМА В МОДУЛЕ КАПЛАНСКОГО — ГИЛЬБЕРТА НАД $L_0$

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Аннотация. В статье изучаются некоторые характеристические свойства самосопряженных частично интегральных операторов типа Фредгольма в модуле Капланского — Гильберта  $L_0\left[L_2\left(\Omega_1\right)\right]$  над  $L_0\left(\Omega_2\right)$ . Используется математический инструментарий из теории модулей Капланского — Гильберта. В модуле Капланского — Гильберта  $L_0\left[L_2\left(\Omega_1\right)\right]$  над  $L_0\left(\Omega_2\right)$  рассматриваются частично интегральные операторы типа Фредгольма  $T_1\left(\Omega_1$  и  $\Omega_2$  — замкнутые ограниченные множества в  $\mathbb{R}^{\nu_1}$  и  $\mathbb{R}^{\nu_2}$ ,  $\nu_1, \nu_2 \in \mathbb{N}$  соответственно). В работе доказано существование  $L_0\left(\Omega_2\right)$ -собственных значений, отличных от нуля для любого самосопряженного частично интегрального оператора типа Фредгольма  $T_1$ ; более того, показано существование конечного или счетного числа вещественных  $L_0\left(\Omega_2\right)$ -собственных значений. В последнем случае, последовательности  $L_0\left(\Omega_2\right)$ -собственных значений порядково сходятся к нулевой функции. Установлена также теорема о разложимости оператора  $T_1$  в ряд по  $\nabla_1$  одномерным операторам.

**Ключевые слова:** частично интегральный оператор, модуль Капланского — Гильберта,  $L_0$ -собственное значение.

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