

GENERALIZED CONNECTION ON $T(T^2M)$

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Abstract. The geometry of some manifolds fibered over a given manifold M is in the first place characterized by the group of allowable coordinate transformations. For the tangent manifold TM these are given by $x^{i'} = x^{i'}(x)y^{i'} = \frac{\partial x^{i'}}{\partial x^i}y^i$, $\text{rank} \left[\frac{\partial x^{i'}}{\partial x^i} \right] = n$, and for the total space of a vector bundle $E \rightarrow M$, we have $x^{i'} = x^{i'}(x)$, $y^{a'} = M_a^{a'}(x)y^a$, $\text{rank}(M_a^{a'}) = m = \text{dimension of type fiber}$.

In the last years R. Miron, Gh. Atanasiu and others examined the $\text{Osc}^k M$ spaces, [10], [11], [12]. Here the case $k = 2$ will be investigated. Instead of $\text{Osc}^2 M$ the notation T^2M will be used ($\text{Osc}^1 M$ coincides with TM). Instead of d -connection used in [10], [11], [12], we consider here the generalized connection and determine its torsion tensor. As a special case the known d -connection is obtained.

1. Adapted basis in $T(T^2M)$. Let T^2M be a $3n$ dimensional C^∞ manifold. A point $u \in T^2M$ in the local charts (U, φ) and (U', φ') has coordinates (x^i, y^i, z^i) and $(x^{i'}, y^{i'}, z^{i'})$ respectively. In $U \cap U'$ the allowed coordinate transformations are given by the equations:

$$(1.1) \quad (a) \ x^{i'} = x^{i'}(x) \quad (b) \ y^{i'} = \frac{\partial x^{i'}}{\partial x^j} y^j \quad (c) \ z^{i'} = \frac{1}{2} \frac{\partial^2 x^{i'}}{\partial x^k \partial x^j} y^k y^j + \frac{\partial x^{i'}}{\partial x^j} z^j.$$

If $\text{rank} \left[\frac{\partial x^{i'}}{\partial x^i} \right] = n$, then the inverse transformation of (1.1) exists:

$$(1.2) \quad (a) \ x^i = x^i(x^{i'}) \quad (b) \ y^i = \frac{\partial x^i}{\partial x^{j'}} y^{j'} \quad (c) \ z^k = \frac{1}{2} \frac{\partial^2 x^k}{\partial x^{i'} \partial x^{j'}} y^{i'} y^{j'} + \frac{\partial x^k}{\partial x^{i'}} z^{i'}.$$

(1.2) (a) and (1.2) (b) are obvious. To obtain (1.2) (c) we start from

$$(1.2) \quad \frac{\partial x^{i'}}{\partial x^h} \frac{\partial x^k}{\partial x^{i'}} = \delta_h^k.$$

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From (1.2) it follows:

$$\frac{\partial^2 x^{i'}}{\partial x^h \partial x^j} \frac{\partial x^k}{\partial x^{i'}} + \frac{\partial x^{i'}}{\partial x^h} \frac{\partial^2 x^k}{\partial x^{i'} \partial x^{j'}} \frac{\partial x^{j'}}{\partial x^j} = 0.$$

The multiplication of the above equation with $y^h y^j$ gives

$$(1.3) \quad \frac{\partial^2 x^{i'}}{\partial x^h \partial x^j} y^h y^j \frac{\partial x^k}{\partial x^{i'}} = - \frac{\partial^2 x^k}{\partial x^{i'} \partial x^{j'}} y^{i'} y^{j'}.$$

If we multiply (1.1) (c) by $\partial x^k / \partial x^{i'}$, use (1.3) and (1.4) we obtain (1.2) (c).

The identity transformation is a special case of (1.1), namely if we put $x^{i'} = x^i$ in (1.1) (a), then $y^{i'} = y^i$, $z^{i'} = z^i$ follows.

If in the local chart (U'', φ'') the point u has coordinates $(x^{i''}, y^{i''}, z^{i''})$, then in $(U'' \cap U', \varphi'')$ (1.1) are valid, if the index i' is substituted by i'' and the indices without $'$ obtain $''$. After some calculation it can be obtained, that the connection between $(x^{i''}, y^{i''}, z^{i''})$ and (x^i, y^i, z^i) in $U \cap U' \cap U''$ is given by (1.1) if the index i' is substituted by i'' .

From the above follows:

THEOREM 1.1. *The transformations of type (1.1) form a group.*

In $T(T^2M)$ the natural bases are:

$$(1.4) \quad \overline{B} = \left\{ \frac{\partial}{\partial x^i}, \frac{\partial}{\partial y^i}, \frac{\partial}{\partial z^i} \right\} \text{ and } \overline{B}' = \left\{ \frac{\partial}{\partial x^{i'}}, \frac{\partial}{\partial y^{i'}}, \frac{\partial}{\partial z^{i'}} \right\}.$$

The bases vectors of \overline{B} and \overline{B}' are connected by:

$$(1.5) \quad \begin{aligned} \frac{\partial}{\partial x^i} &= \frac{\partial x^{i'}}{\partial x^i} \frac{\partial}{\partial x^{i'}} + \frac{\partial y^{i'}}{\partial x^i} \frac{\partial}{\partial y^{i'}} + \frac{\partial z^{i'}}{\partial x^i} \frac{\partial}{\partial z^{i'}}, \\ \frac{\partial}{\partial y^i} &= \frac{\partial y^{i'}}{\partial y^i} \frac{\partial}{\partial y^{i'}} + \frac{\partial z^{i'}}{\partial y^i} \frac{\partial}{\partial z^{i'}}, \\ \frac{\partial}{\partial z^i} &= \frac{\partial z^{i'}}{\partial z^i} \frac{\partial}{\partial z^{i'}}. \end{aligned}$$

From (1.1) (c) it follows

$$(1.6) \quad \frac{\partial}{\partial z^i} = \frac{\partial x^{i'}}{\partial x^i} \frac{\partial}{\partial z^{i'}}.$$

Let us introduce the notation

$$(1.7) \quad \frac{\delta}{\delta y^i} = \frac{\partial}{\partial y^i} - \mathcal{H}_i^j(x, y) \frac{\partial}{\partial z^i}.$$

PROPOSITION 1.1. $\frac{\delta}{\delta y^i}$ defined by (1.8) is transformed as tensor, i.e.

$$(1.8) \quad \frac{\delta}{\delta y^i} = \frac{\partial x^{i'}}{\partial x^i} \frac{\delta}{\delta y^{i'}}$$

if $\mathcal{H}_j^i(x, y)$ is transformed in the following way:

$$(1.9) \quad \mathcal{H}_{i'}^{j'} = \mathcal{H}_i^j \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^{j'}}{\partial x^j} - \frac{\partial^2 x^{j'}}{\partial x^i \partial x^j} y^j \frac{\partial x^i}{\partial x^{i'}}.$$

The proof is obtained by direct calculation.

From (1.9) and (1.10) we obtain the first connection coefficient of Berwald type:

$$(1.10) \quad \mathcal{H}_{i'k'}^{j'} = \frac{\partial \mathcal{H}_{i'}^{j'}}{\partial y^{k'}} = \frac{\partial \mathcal{H}_i^j}{\partial y^k} \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^{j'}}{\partial x^j} \frac{\partial x^k}{\partial x^{k'}} - \frac{\partial^2 x^{j'}}{\partial x^i \partial x^k} \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^k}{\partial x^{k'}}.$$

Let us introduce the notation:

$$(1.11) \quad \frac{\delta}{\delta x^i} = \frac{\partial}{\partial x^i} - \mathcal{M}_i^j(x, y) \frac{\partial}{\partial y^j} - \mathcal{N}_i^j(x, y, z) \frac{\partial}{\partial z^j}.$$

PROPOSITION 1.2. $\frac{\delta}{\delta x^i}$ defined by (1.12) is transformed in the form:

$$(1.12) \quad \frac{\delta}{\delta x^i} = \frac{\partial x^{i'}}{\partial x^i} \frac{\delta}{\delta x^{i'}}$$

if $\mathcal{M}_i^j(x, y)$ and $\mathcal{N}_i^j(x, y, z)$ are transformed in the following way:

$$(1.13) \quad \mathcal{M}_{i'}^{j'} = \mathcal{M}_i^j \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^{j'}}{\partial x^j} - \frac{\partial^2 x^{j'}}{\partial x^i \partial x^j} y^j \frac{\partial x^i}{\partial x^{i'}}$$

$$(1.14) \quad \begin{aligned} \mathcal{N}_{i'}^{j'} &= \mathcal{N}_i^j \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^{j'}}{\partial x^j} - \frac{\partial^2 x^{j'}}{\partial x^i \partial x^j} z^j \frac{\partial x^i}{\partial x^{i'}} \\ &\quad - \frac{1}{2} \frac{\partial^3 x^{j'}}{\partial x^i \partial x^h \partial x^k} y^h y^k \frac{\partial x^i}{\partial x^{i'}} + \frac{\partial^2 x^{j'}}{\partial x^k \partial x^j} y^k \mathcal{M}_h^j \frac{\partial x^h}{\partial x^{i'}}. \end{aligned}$$

Remark. \mathcal{M}_i^j and \mathcal{N}_i^j used here in [10], [11], [12] are denoted by (1) \mathcal{N}_i^j and (2) \mathcal{N}_i^j respectively.

Proof. If we add the following equations (which follow from (1.6), (1.1) and (1.3)):

$$\begin{aligned} \frac{\partial}{\partial x^i} &= \left(\frac{\partial x^{i'}}{\partial x^i} \right) \frac{\partial}{\partial x^{i'}} + \left(\frac{\partial^2 x^{j'}}{\partial x^k \partial x^j} y^j \frac{\partial x^k}{\partial x^{i'}} \frac{\partial x^{i'}}{\partial x^i} \right) \frac{\partial}{\partial y^{j'}} \\ &\quad + \left(\frac{1}{2} \frac{\partial^3 x^{j'}}{\partial x^h \partial x^k \partial x^j} y^k y^j \frac{\partial x^h}{\partial x^{i'}} \frac{\partial x^{i'}}{\partial x^i} + \frac{\partial^2 x^{j'}}{\partial x^k \partial x^j} z^j \frac{\partial x^k}{\partial x^{i'}} \frac{\partial x^{i'}}{\partial x^i} \right) \frac{\partial}{\partial z^{j'}} \\ -\mathcal{M}_i^j \frac{\partial}{\partial y^j} &= -\mathcal{M}_h^j \frac{\partial x^h}{\partial x^{i'}} \frac{\partial x^{i'}}{\partial x^i} \left(\frac{\partial x^{j'}}{\partial x^j} \frac{\partial}{\partial y^{j'}} + \frac{\partial^2 x^{j'}}{\partial x^k \partial x^j} y^k \frac{\partial}{\partial z^{j'}} \right), \\ -\mathcal{N}_i^j \frac{\partial}{\partial z^j} &= -\mathcal{N}_h^j \frac{\partial x^{j'}}{\partial x^j} \frac{\partial x^h}{\partial x^{i'}} \frac{\partial x^{i'}}{\partial x^i} \frac{\partial}{\partial z^{j'}} , \end{aligned}$$

we obtain $\frac{\delta}{\delta x^i} = \frac{\partial x^{i'}}{\partial x^i} \left(\frac{\partial}{\partial x^{i'}} - \mathcal{M}_{i'}^{j'} \frac{\partial}{\partial y^{j'}} - \mathcal{N}_{i'}^{j'} \frac{\partial}{\partial z^{j'}} \right) = \frac{\partial x^{i'}}{\partial x^i} \frac{\delta}{\delta x^{i'}}$, where $\mathcal{M}_{i'}^{j'}$ is given by (1.14) and $\mathcal{N}_{i'}^{j'}$ by (1.15).

The basis

$$(1.15) \quad B = \left\{ \frac{\delta}{\delta x^i}, \frac{\delta}{\delta y^i}, \frac{\partial}{\partial z^i} \right\}$$

is called *adapted basis* for $T(T^2M)$. Its elements are transformed by (1.7), (1.9) and (1.13).

From (1.14) and (1.15) we obtain two other connection coefficients of Berwald type, namely

$$(1.16) \quad \mathcal{M}_{i'k'}^{j'} = \frac{\partial \mathcal{M}_{i'}^{j'}}{\partial x^{k'}} = \mathcal{M}_{ik}^j \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^k}{\partial x^{k'}} \frac{\partial x^{j'}}{\partial x^j} - \frac{\partial^2 x^{j'}}{\partial x^i \partial x^k} \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^k}{\partial x^{k'}}.$$

$$(1.17) \quad \mathcal{N}_{i'k'}^{j'} = \frac{\partial \mathcal{N}_{i'}^{j'}}{\partial z^{k'}} = \mathcal{N}_{ik}^j \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^k}{\partial x^{k'}} \frac{\partial x^{j'}}{\partial x^j} - \frac{\partial^2 x^{j'}}{\partial x^i \partial x^k} \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^k}{\partial x^{k'}}.$$

2. The adapted basis in $T^*(T^2M)$. The natural basis of $T^*(T^2M)$ is

$$(2.1) \quad \bar{B}^* = \{dx^i, dy^i, dz^i\},$$

where the following relations are valid with respect to (1.1):

$$(2.2) \quad (a) \quad dx^{i'} = \frac{\partial x^{i'}}{\partial x^i} dx^i \quad (b) \quad dy^{i'} = \frac{\partial^2 x^{i'}}{\partial x^j \partial x^i} y^j dx^j + \frac{\partial x^{i'}}{\partial x^i} dy^i$$

$$(c) \quad dz^{i'} = \left(\frac{1}{2} \frac{\partial^3 x^{i'}}{\partial x^k \partial x^h \partial x^j} y^k y^h + \frac{\partial^2 x^{i'}}{\partial x^i \partial x^j} z^i \right) dx^j + \left(\frac{\partial^2 x^{i'}}{\partial x^k \partial x^j} y^k \right) dy^j + \left(\frac{\partial x^{i'}}{\partial x^j} \right) dz^j.$$

From (2.2) (b) and (2.2) (c) it is obvious, that the transformations of $dy^{i'}$ and $dz^{i'}$ are not tensorial. If we put

$$(2.3) \quad \delta y^i = dy^i + \mathcal{H}_j^i dx^j$$

then from (1.10) and (2.1) we get

$$(2.4) \quad \delta y^{i'} = \frac{\partial x^{i'}}{\partial x^i} \delta y^i.$$

Let us introduce the notation

$$(2.5) \quad \delta z^i = dz^i + \mathcal{M}_j^i(x, y) dy^j + \mathcal{N}_j^i(x, y, z) dx^j + \mathcal{G}_j^i(x, y) dx^j,$$

where the transformation laws of \mathcal{M}_j^i and \mathcal{N}_j^i are prescribed by (1.14) and (1.15).

PROPOSITION 2.1. *If $\mathcal{G}_j^i(x, y)$ has the following law of transformation (with respect to (1.1)):*

$$(2.6) \quad \mathcal{G}_j^i = \mathcal{G}_{j'}^{i'} \frac{\partial x^i}{\partial x^{i'}} \frac{\partial x^{j'}}{\partial x^j} + \mathcal{M}_{j'}^{i'} \frac{\partial^2 x^{j'}}{\partial x^k \partial x^j} y^k \frac{\partial x^i}{\partial x^{i'}} + \mathcal{M}_j^k \frac{\partial^2 x^{j'}}{\partial x^k \partial x^h} y^h \frac{\partial x^i}{\partial x^{j'}},$$

then δz^i defined by (2.5) satisfies the equation

$$(2.7) \quad \delta z^{j'} = \frac{\partial x^{j'}}{\partial x^j} \delta z^j$$

Proof. The substitution of (2.2), (1.14) and (1.15) into $\delta z^{j'} = dz^{j'} + \mathcal{M}_{i'}^{j'} dy^{i'} + \mathcal{N}_{i'}^{j'} dx^{i'}$ gives (2.7) if (2.6) is true.

The adapted basis of $T^*(T^2F)$ is

$$(2.8) \quad B^* = \{dx^i, \delta y^i, \delta z^i\},$$

where the elements of B^* defined by (2.3) and (2.5) satisfy the transformation law prescribed by (2.2) (a), (2.4) and (2.7).

THEOREM 2.1. *The bases B (1.16) and B^* (2.8) of $T(T^2M)$ and $T^*(T^2M)$ respectively are dual to each other if*

$$(2.9) \quad \mathcal{H}_j^i(x, y) = \mathcal{M}_j^i(x, y),$$

$$(2.10) \quad \mathcal{G}_j^i(x, y) = \mathcal{M}_k^i(x, y) \mathcal{M}_j^k(x, y)$$

and if \overline{B} (1.6) and \overline{B}^* (2.1) are dual to each other.

Proof. By direct calculation using (1.8), (1.12), (2.3) and (2.5) we obtain:

$$(2.11) \quad \begin{aligned} \langle dx^j, \frac{\delta}{\delta x^i} \rangle &= \delta_i^j, & \langle dx^j, \frac{\delta}{\delta y^i} \rangle &= 0, & \langle dx^j, \frac{\partial}{\partial z^i} \rangle &= 0, \\ \langle \delta y^j, \frac{\delta}{\delta x^i} \rangle &= \mathcal{H}_i^j - \mathcal{M}_i^j, & \langle \delta y^j, \frac{\delta}{\delta y^i} \rangle &= \delta_i^j, & \langle \delta y^j, \frac{\partial}{\partial z^i} \rangle &= 0, \\ \langle \delta z^j, \frac{\delta}{\delta x^i} \rangle &= \mathcal{G}_i^j - \mathcal{M}_k^j \mathcal{M}_i^k, & \langle \delta z^j, \frac{\delta}{\delta y^i} \rangle &= -\mathcal{H}_i^j + \mathcal{M}_i^j, & \langle \delta z^j, \frac{\partial}{\partial z^i} \rangle &= \delta_i^j. \end{aligned}$$

The duality follows from (2.11) and (2.9).

PROPOSITION 2.2. *$\mathcal{G}_i^j(x, y)$ defined by (2.10) satisfies (2.6).*

Proof. Using (1.4) and (1.14), it can be proved that

$$(2.12) \quad \mathcal{M}_r^s = \mathcal{M}_{i'}^{j'} \frac{\partial x^{i'}}{\partial x^r} \frac{\partial x^s}{\partial x^{j'}} - \frac{\partial^2 x^s}{\partial x^{j'} \partial x^{i'}} y^{i'} \frac{\partial x^{j'}}{\partial x^r}.$$

From (2.12) and (2.10) after some calculation we obtain (2.6).

Remark. It is important that the bases B and B^* be dual to each other, because if they are not, then the contraction of tensors doesn't result tensors. If \overline{B} and \overline{B}^* ((1.6) and (2.1)) are dual to each other, it doesn't follow that \overline{B}' and \overline{B}'^* are dual.

Now we have:

THEOREM 2.2. *The bases $B(\mathcal{M}, \mathcal{N}) = \left\{ \frac{\delta}{\delta x^i}, \frac{\delta}{\delta y^i}, \frac{\partial}{\partial z^i} \right\}$ and $B^*(\mathcal{M}, \mathcal{N}) = \{dx^i, \delta y^i, \delta z^i\}$, where their elements are given by*

$$(2.13) \quad \begin{aligned} \frac{\delta}{\delta x^i} &= \frac{\partial}{\partial x^i} - \mathcal{M}_i^j \frac{\partial}{\partial y^j} - \mathcal{N}_i^j \frac{\partial}{\partial z^j}, & \frac{\delta}{\delta y^i} &= \frac{\partial}{\partial y^i} - \mathcal{M}_i^j \frac{\partial}{\partial z^j}, \\ \delta y^i &= dy^i + \mathcal{M}_j^i dx^j, & \delta z^i &= dz^i + \mathcal{M}_j^i \delta y^j + \mathcal{N}_j^i dx^j \end{aligned}$$

are adapted basis for $T(T^2M)$ and $T^*(T^2M)$ respectively, dual to each other, and they satisfy the law of transformation:

$$(2.14) \quad \begin{aligned} \frac{\delta}{\delta x^i} &= \frac{\partial x^{i'}}{\partial x^i} \frac{\delta}{\delta x^{i'}} + \frac{\delta}{\delta y^{i'}} \frac{\partial y^{i'}}{\partial x^i} + \frac{\partial}{\partial z^{i'}} \frac{\partial z^{i'}}{\partial x^i}, \\ dx^i &= \frac{\partial x^i}{\partial x^{i'}} dx^{i'}, & \delta y^i &= \frac{\partial y^i}{\partial y^{i'}} \delta y^{i'}, & \delta z^i &= \frac{\partial z^i}{\partial z^{i'}} \delta z^{i'}. \end{aligned}$$

It must be noted that there exist as many adapted bases as many functions $\mathcal{M}_i^j(x, y)$ and $\mathcal{N}_i^j(x, y, z)$ can be found, satisfying (1.14) and (1.15) respectively.

If we denote by T_H, T_{V_1}, T_{V_2} the subspaces of $T(T^2M)$ spanned by $\left\{ \frac{\delta}{\delta x^i} \right\}$, $\left\{ \frac{\delta}{\delta y^i} \right\}$, $\left\{ \frac{\partial}{\partial z^i} \right\}$, and by $T_H^*, T_{V_1}^*, T_{V_2}^*$ the subspaces of $T^*(T^2M)$ spanned by $\{dx^i\}$, $\{\delta y^i\}$, $\{\delta z^i\}$ respectively, then

$$T(T^2M) = T_H \oplus T_{V_1} \oplus T_{V_2}, \quad T^*(T^2M) = T_H^* \oplus T_{V_1}^* \oplus T_{V_2}^*.$$

For the further examinations it is useful to introduce different kinds of indices. Indices $i, j, h, k, l = \overline{1, n}$ will be used in T_H and T_H^* , $a, b, c, d, e, f = \overline{n+1, 2n}$ in T_{V_1} and $T_{V_1}^*$, $p, q, r, s, t = \overline{2n+1, 3n}$ in T_{V_2} and $T_{V_2}^*$. The Greek letters as indices will take values from 1 to $3n$. Using this notation the adapted bases have the form:

$$(2.15) \quad B = \left\{ \frac{\delta}{\delta x^i}, \frac{\delta}{\delta y^a}, \frac{\partial}{\partial z^p} \right\} = \{\delta_\alpha\}, \quad B^* = \{dx^j, \delta y^b, \delta z^q\} = \{\delta^\beta\},$$

where

$$(2.16) \quad \begin{aligned} \frac{\delta}{\delta x^i} &= \frac{\partial}{\partial x^i} - \mathcal{M}_i^b \frac{\partial}{\partial y^b} - \mathcal{N}_i^q \frac{\partial}{\partial z^q}, & \frac{\delta}{\delta y^a} &= \frac{\partial}{\partial y^a} - \mathcal{M}_a^q \frac{\partial}{\partial z^q}, \\ \delta y^a &= dy^a + \mathcal{M}_i^a dx^i, & \delta z^q &= dz^q + \mathcal{M}_a^q \delta y^a + \mathcal{N}_j^q dx^j. \end{aligned}$$

If $i = a \pmod{n}$ and $j = b = q \pmod{n}$, then $\mathcal{M}_i^j = \mathcal{M}_i^b = \mathcal{M}_a^q$, $\mathcal{N}_i^q = \mathcal{N}_i^j$, in (2.16).

Some tensor field T expressed in the bases B and B^* ((2.15)) has the form:

$$T = T_{\dots j \dots a \dots r \dots} \frac{\delta}{\delta x^i} \otimes dx^j \dots \frac{\delta}{\delta y^a} \otimes \delta y^b \dots \frac{\partial}{\partial z^r} \otimes \delta z^s \dots$$

The components of the tensor T , with respect to the coordinate transformations (1.1) are transformed in the following way:

$$T_{\dots j' \dots a' \dots r' \dots} = T_{\dots j \dots a \dots r \dots} \frac{\partial x^{i'}}{\partial x^i} \frac{\partial x^j}{\partial x^{j'}} \dots \frac{\partial y^{a'}}{\partial y^a} \frac{\partial y^b}{\partial y^{b'}} \dots \frac{\partial z^{r'}}{\partial z^r} \frac{\partial z^s}{\partial z^{s'}} \dots$$

For some vector field $X \in T(T^2M)$ and some 1-form $\omega \in T^*(T^2M)$ we have:

$$(2.17) \quad \begin{aligned} X &= X^i \frac{\delta}{\delta x^i} + X^a \frac{\delta}{\delta y^a} + X^p \frac{\partial}{\partial z^p} = X^\alpha \delta_\alpha, \\ \omega &= \omega_j dx^j + \omega_b \delta y^b + \omega_q \delta z^q = \omega_\beta \delta^\beta. \end{aligned}$$

With respect to (1.1) the coordinates of X and ω transform in the following way:

$$\begin{aligned} X^{i'} &= X^i \frac{\partial x^{i'}}{\partial x^i}, & X^{a'} &= X^a \frac{\partial y^{a'}}{\partial y^a}, & X^{p'} &= X^p \frac{\partial z^{p'}}{\partial z^p}, \\ \omega_{j'} &= \omega_j \frac{\partial x^j}{\partial x^{j'}}, & \omega_{b'} &= \omega_b \frac{\partial y^b}{\partial y^{b'}}, & \omega_{q'} &= \omega_q \frac{\partial z^q}{\partial z^{q'}}, \end{aligned}$$

because for $i = a = p \pmod{n}$ we have:

$$(2.18) \quad \frac{\partial x^{i'}}{\partial x^i} = \frac{\partial y^{a'}}{\partial y^a} = \frac{\partial z^{p'}}{\partial z^p}.$$

3. Generalized covariant derivatives. The generalized connection in Lagrange and Hamilton spaces was studied among others in [2]–[6]. In $T(T^2M)$ it is introduced in the following way. Let $\nabla : T(T^2M) \times T(T^2M) \rightarrow T(T^2M)$ (\times is the Descartes product) be a linear connection, such that $\nabla : (X, Y) \rightarrow \nabla_X Y \in T(T^2M)$, $\forall X, Y \in T(T^2M)$. The operator ∇ is called generalized connection. It is called d -connection if $\nabla_X Y$ is in T_H , T_{V_1} or T_{V_2} if Y is in T_H , T_{V_1} or T_{V_2} , respectively, $\forall X \in T(T^2M)$. It has been studied by many authors, mostly romanian geometers.

We shall not make that restriction on ∇ here. In the following we shall use the abbreviations: $\delta_k = \frac{\delta}{\delta x^k}$, $\delta_a = \frac{\delta}{\delta y^a}$, $\partial_k = \frac{\partial}{\partial x^k}$, $\partial_a = \frac{\partial}{\partial y^a}$, $\partial_p = \frac{\partial}{\partial z^p}$.

Definition 3.1. The generalized connection ∇ is defined by

$$(3.1) \quad \begin{aligned} \nabla_{\delta_i} \delta_j &= \underline{F}_{j_i}^k \delta_k + F_{j_i}^c \delta_c + F_{j_i}^r \partial_r, & \nabla_{\delta_i} \delta_b &= F_{b_i}^k \delta_k + \underline{F}_{b_i}^c \delta_c + F_{b_i}^r \partial_r \\ \nabla_{\delta_i} \partial_q &= F_{q_i}^k \delta_k + F_{q_i}^c \delta_c + \underline{F}_{q_i}^r \partial_r, & \nabla_{\delta_a} \delta_j &= \underline{C}_{j_a}^k \delta_k + C_{j_a}^c \delta_c + C_{j_a}^r \partial_r, \\ \nabla_{\delta_a} \delta_b &= C_{b_a}^k \delta_k + \underline{C}_{b_a}^c \delta_c + C_{b_a}^r \partial_r, & \nabla_{\delta_a} \partial_q &= C_{q_a}^k \delta_k + C_{q_a}^c \delta_c + \underline{C}_{q_a}^r \partial_r, \\ \nabla_{\partial_p} \delta_j &= \underline{L}_{j_p}^k \delta_k + L_{j_p}^c \delta_c + L_{j_p}^r \partial_r, & \nabla_{\partial_p} \delta_b &= L_{b_p}^k \delta_k + \underline{L}_{b_p}^c \delta_c + L_{b_p}^r \partial_r, \\ \nabla_{\partial_p} \partial_q &= L_{q_p}^k \delta_k + L_{q_p}^c \delta_c + \underline{L}_{q_p}^r \partial_r. \end{aligned}$$

The d -connection is defined if in (3.1) all terms on the right-hand side vanish, except the underlined ones.

For the vector field X defined by (2.17) we have

$$\begin{aligned} &\nabla_{\delta_i} (X^j \delta_j + X^b \delta_b + X^q \partial_q) \\ &= (\delta_i X^j) \delta_j + X^j \nabla_{\delta_i} \delta_j + (\delta_i X^b) \delta_b + X^b \nabla_{\delta_i} \delta_b + (\delta_i X^q) \partial_q + X^q \nabla_{\delta_i} \partial_q \\ &= (\delta_i X^k + F_{j_i}^k X^j + F_{b_i}^k X^b + F_{q_i}^k X^q) \delta_k + (\delta_i X^c + F_{j_i}^c X^j + F_{b_i}^c X^b + F_{q_i}^c X^q) \delta_c + \\ &\quad (\delta_i X^r + F_{j_i}^r X^j + F_{b_i}^r X^b + F_{q_i}^r X^q) \partial_r \end{aligned}$$

From the above equation it follows

$$(3.2) \quad \nabla_{\delta_i} X = X_{|i}^k \delta_k + X_{|i}^c \delta_c + X_{|i}^r \partial_r,$$

where

$$(3.3) \quad X_{|i}^x = \delta_i X^x + F_{j_i}^x X^j + F_{b_i}^x X^b + F_{q_i}^x X^q, \quad x \in \{k, c, r\}$$

or shorter

$$(3.4) \quad X_{|i}^x = \delta_i X^x + F_{\alpha_i}^x X^\alpha, \quad x \in \{k, c, r\}.$$

The summation over α is the sum of summations over j , b and q as is written in (3.3). The sign $|i$ is the covariant derivative in direction of the basis vector δ_i .

The covariant derivative of X in the direction of δ_a has the form:

$$(3.5) \quad \nabla_{\delta_a} X = X^k|_a \delta_k + X^c|_a \delta_c + X^r|_a \partial_r,$$

where

$$(3.6) \quad X^x|_a = \delta_a X^x + C_{j_a}^x X^j + C_{b_a}^x X^b + C_{q_a}^x X^q, \quad x \in \{k, c, r\},$$

or shorter

$$(3.7) \quad X^x|_a = \delta_a X^x + C_{\alpha_a}^x X^\alpha.$$

The covariant derivative of the vector field X in the direction of ∂_p is given by

$$(3.8) \quad \nabla_{\partial_p} X = X^k \parallel_p \delta_k + X^c \parallel_p \delta_c + X^r \parallel_p \partial_r,$$

where

$$(3.9) \quad X^x \parallel_p = \partial_p X^x + L_{j_p}^x X^j + L_{b_p}^x X^b + L_{q_p}^x X^q, \quad x \in \{k, c, r\},$$

or abbreviated

$$(3.10) \quad X^x \parallel_p = \partial_p X^x + L_{\alpha_p}^x X^\alpha.$$

In (3.7) and (3.10) the summation over α is the sum of summations over j , b and q (as in (3.4)).

THEOREM 3.1. *If X and Y are vector fields in $T(T^2M)$, ∇ the generalized connection defined by (3.1), then the following equation is valid:*

$$(3.11) \quad \nabla_Y X = (X^k \parallel_i Y^i + X^k \parallel_a Y^a + X^k \parallel_p Y^p) \delta_k + (X^c \parallel_i Y^i + X^c \parallel_a Y^a + X^c \parallel_p Y^p) \delta_c + (X^r \parallel_i Y^i + X^r \parallel_a Y^a + X^r \parallel_p Y^p) \partial_r.$$

Proof. The proof follows from (3.2)–(3.10) and the bilinearity of ∇ .

The equation (3.11) can be written in the abbreviated form as follows

$$(3.12) \quad \nabla_Y X = X^{\alpha} \parallel_{\beta} Y^{\beta} \delta_{\alpha},$$

$$(3.13) \quad X^{\alpha} \parallel_{\beta} = \delta_{\beta} X^{\alpha} + \Gamma_{\gamma\beta}^{\alpha} X^{\gamma}.$$

If $\beta = i$, then $\Gamma = F$; if $\beta = a$, then $\Gamma = C$; if $\beta = p$ then $\Gamma = L$.

THEOREM 3.2. *All covariant derivatives $X^{\alpha} \parallel_i$, $X^{\alpha} \parallel_a$, $X^{\alpha} \parallel_p$ ($\alpha = k$, or $\alpha = c$ or $\alpha = r$) from (3.11) are transformed as tensors with respect to (1.1) if all connection coefficients from (3.1) are transformed as tensors, except the following, which have the form*

$$(3.14) \quad \begin{aligned} F_{j_i}^k &= F_{j' i'}^{k'} \frac{\partial x^{i'}}{\partial x^i} \frac{\partial x^k}{\partial x^{k'}} \frac{\partial x^{j'}}{\partial x^j} + \frac{\partial^2 x^{k'}}{\partial x^i \partial x^j} \frac{\partial x^k}{\partial x^{k'}} \\ F_{b_i}^c &= F_{b' i'}^{c'} \frac{\partial x^{i'}}{\partial x^i} \frac{\partial y^{b'}}{\partial y^b} \frac{\partial y^c}{\partial y^{c'}} + \frac{\partial^2 y^{c'}}{\partial x^i \partial y^b} \frac{\partial y^c}{\partial y^{c'}} \\ F_{q_i}^r &= F_{q' i'}^{r'} \frac{\partial x^{i'}}{\partial x^i} \frac{\partial z^{q'}}{\partial z^q} \frac{\partial z^r}{\partial z^{r'}} + \frac{\partial^2 z^{r'}}{\partial x^i \partial z^q} \frac{\partial z^r}{\partial z^{r'}}. \end{aligned}$$

The proof is obtained by direct calculation.

Remark. From (2.17) it follows that $F_{j i}^k = F_{b i}^c = F_{q i}^r$ if $k = c = r \pmod{n}$, $j = b = q \pmod{n}$. The connection coefficients $\mathcal{H}_{j k}^i$, $\mathcal{M}_{i k}^j$ and $\mathcal{N}_{i k}^j$ defined by (1.11), (1.17) and (1.18) respectively, satisfy the transformation laws prescribed by (3.14).

4. The torsion tensor of the generalized connection. The torsion tensor $T(X, Y)$ is defined in the usual way by:

$$T(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y].$$

THEOREM 4.1. *The torsion tensor for the generalized connection has the form $T(X, Y) = T^k \delta_k + T^c \delta_c + T^r \partial_r$, where*

$$\begin{aligned} T^k &= T_{j i}^k Y^j X^i + T_{j b}^k Y^j X^b + T_{j q}^k Y^j X^q \\ &\quad + T_{b i}^k Y^b X^i + T_{b a}^k Y^b X^a + T_{b q}^k Y^b X^q \\ &\quad + T_{p i}^k Y^p X^i + T_{p b}^k Y^p X^b + T_{p q}^k Y^p X^q, \\ T_{j i}^k &= F_{j i}^k - F_{i j}^k, \quad T_{j b}^k = C_{j b}^k - F_{b j}^k, \quad T_{j q}^k = L_{j q}^k - F_{q j}^k, \\ T_{b i}^k &= F_{b i}^k - C_{i b}^k, \quad T_{b a}^k = C_{b a}^k - C_{a b}^k, \quad T_{b q}^k = L_{b q}^k - C_{q b}^k, \\ T_{p i}^k &= F_{p i}^k - L_{i p}^k, \quad T_{p b}^k = C_{p b}^k - L_{b p}^k, \quad T_{p q}^k = L_{p q}^k - L_{q p}^k, \\ T^c &= T_{j i}^c Y^j X^i + T_{j b}^c Y^j X^b + T_{j r}^c Y^j X^r + T_{b i}^c Y^b X^i + T_{a b}^c Y^a X^b + \\ &\quad T_{b r}^c Y^b X^r + T_{q i}^c Y^q X^i + T_{q b}^c Y^q X^b + T_{q r}^c Y^q X^r, \\ T_{j i}^c &= F_{j i}^c - F_{i j}^c - K_{i j}^c, \quad T_{a b}^c = C_{a b}^c - C_{b a}^c, \quad T_{q i}^c = F_{q i}^c - L_{i q}^c, \\ T_{j b}^c &= C_{j b}^c - F_{b j}^c + K_{j b}^c, \quad T_{b r}^c = L_{b r}^c - C_{r b}^c, \quad T_{q b}^c = C_{q b}^c - L_{b q}^c, \\ T_{b i}^c &= F_{b i}^c - C_{i b}^c - K_{i b}^c, \quad T_{j r}^c = L_{j r}^c - F_{r j}^c, \quad T_{q r}^c = L_{q r}^c - L_{r q}^c, \\ T^r &= T_{j i}^r Y^j X^i + T_{j b}^r Y^j X^b + T_{j q}^r Y^j X^q + T_{b i}^r Y^b X^i + T_{b a}^r Y^b X^a + \\ &\quad T_{b q}^r Y^b X^q + T_{p i}^r Y^p X^i + T_{p b}^r Y^p X^b + T_{p q}^r Y^p X^q \\ T_{j i}^r &= F_{j i}^r - F_{i j}^r - K_{i j}^r, \quad T_{j b}^r = C_{j b}^r - F_{b j}^r + K_{j b}^r, \quad T_{b q}^r = L_{b q}^r - C_{q b}^r, \\ T_{j q}^r &= L_{j q}^r - F_{q j}^r + K_{j q}^r, \quad T_{b i}^r = F_{b i}^r - C_{i b}^r - K_{i b}^r, \quad T_{p b}^r = C_{p b}^r - L_{b p}^r, \\ T_{b a}^r &= C_{b a}^r - C_{a b}^r - K_{a b}^r, \quad T_{p i}^r = F_{p i}^r - L_{i p}^r - K_{i p}^r, \quad T_{p q}^r = L_{p q}^r - L_{q p}^r. \end{aligned}$$

Proof. The proof is obtained by direct calculation using (3.12), (3.13) and the relations:

$$\begin{aligned} [X, Y] &= [X(Y^j) - Y(X^j)]\delta_j + [X(Y^b) - Y(X^b)]\delta_b + \\ &\quad [X(Y^q) - Y(X^q)]\partial_q + A + B + C + D, \end{aligned}$$

where

$$\begin{aligned} A &= X^i Y^j (\delta_i \delta_j - \delta_j \delta_i), \quad B = (X^i Y^b - Y^i X^b) (\delta_i \delta_b - \delta_b \delta_i) \\ C &= (X^i Y^q - Y^i X^q) (\delta_i \partial_q - \partial_q \delta_i), \quad D = X^a Y^b (\delta_a \delta_b - \delta_b \delta_a). \end{aligned}$$

Using (2.16) we obtain:

$$A = X^i Y^j (K_{ij}^c \partial_c + \bar{K}_{ij}^r \partial_r) = X^i Y^j (K_{ij}^c \delta_c + K_{ij}^r \partial_r),$$

where

$$\begin{aligned} K_{ij}^c &= \delta_j \mathcal{M}_i^c - \delta_i \mathcal{M}_j^c, & K_{ij}^r &= \bar{K}_{ij}^r + \mathcal{M}_c^r K_{ij}^c, & \bar{K}_{ij}^r &= \delta_j \mathcal{N}_i^r - \delta_i \mathcal{N}_j^r. \\ B &= (X^i Y^b - Y^i X^b) (K_{ib}^c \delta_c + K_{ib}^r \partial_r), \\ K_{ib}^c &= \delta_b \mathcal{M}_i^c, & K_{ib}^r &= \bar{K}_{ib}^r + \mathcal{M}_c^r K_{ib}^c, & \bar{K}_{ib}^r &= \delta_b \mathcal{N}_i^r - \delta_i \mathcal{M}_b^r. \\ C &= (X^i Y^a - Y^i X^a) K_{iq}^r \partial_r, & K_{iq}^r &= \frac{\partial \mathcal{N}_i^r}{\partial z^q}. \\ D &= X^a Y^b K_{ab}^r \partial_r, & K_{ab}^r &= \delta_b \mathcal{M}_a^r - \delta_a \mathcal{M}_b^r. \end{aligned}$$

Remark. As $\mathcal{M} = \mathcal{M}(x, y)$, in all above formulae $\partial_r \mathcal{M} = 0$.

5. Special cases. As mentioned in Definition 3.1 the special connection $\bar{\nabla}$ (the so called d -connection) is obtained, if in (3.1) only the underlined terms are left. More precisely:

Definition 5.1 The d -connection $\bar{\nabla}$ is defined by:

$$(5.1) \quad \begin{aligned} \bar{\nabla}_{\delta_i} \delta_j &= \bar{F}_{ji}^k \delta_k, & \bar{\nabla}_{\delta_i} \delta_b &= \bar{F}_{bi}^c \delta_c, & \bar{\nabla}_{\delta_i} \partial_q &= \bar{F}_{qi}^r \partial_r \\ \bar{\nabla}_{\delta_a} \delta_j &= \bar{C}_{ja}^k \delta_k, & \bar{\nabla}_{\delta_a} \delta_b &= \bar{C}_{ba}^c \delta_c, & \bar{\nabla}_{\delta_a} \partial_q &= \bar{C}_{qa}^r \partial_r \\ \bar{\nabla}_{\partial_p} \delta_j &= \bar{L}_{jp}^k \delta_k, & \bar{\nabla}_{\partial_p} \delta_b &= \bar{L}_{bp}^c \delta_c, & \bar{\nabla}_{\partial_p} \partial_q &= \bar{L}_{qp}^r \partial_r. \end{aligned}$$

From (5.1) the following property of d -connection is obvious:

$$\bar{\nabla}_X : T_H \rightarrow T_H, \quad \bar{\nabla}_X : T_{V_1} \rightarrow T_{V_1}, \quad \bar{\nabla}_X : T_{V_2} \rightarrow T_{V_2}$$

for any vector field X from $T(T^2M)$.

THEOREM 5.1. *If X and Y are vector fields in $T(T^2M)$ expressed in the basis B (2.15), then*

$$\begin{aligned} \bar{\nabla}_X Y &= Y_{\bar{i}}^k X^i + Y^k \bar{\lceil}_a X^a + Y^k \bar{\lceil}_p X^p \delta_k + (Y_{\bar{i}}^c X^i + Y^c \bar{\lceil}_a X^a + Y^c \bar{\lceil}_p X^p) \delta_c \\ &\quad + (Y_{\bar{i}}^r X^i + Y^r \bar{\lceil}_a X^a + Y^r \bar{\lceil}_p X^p) \partial_r, \\ Y_{\bar{i}}^x &= \delta_i Y^x + \bar{F}_{yi}^x Y^y, & Y^x \bar{\lceil}_a &= \partial_a Y^x + \bar{C}_{ya}^x Y^y, & Y^x \bar{\lceil}_p &= \partial_p Y^x + \bar{L}_{yp}^x Y^y, \end{aligned}$$

where either $x = k, y = j$, or $x = c, y = b$, or $x = r, y = q$.

THEOREM 5.2. *The connection coefficients $\bar{C}_{ja}^k, \bar{C}_{ba}^c, \bar{C}_{qa}^r, \bar{L}_{jp}^k, \bar{L}_{bp}^c$ and \bar{L}_{qp}^r with respect to (1.1) are transformed as tensors. The transformation laws for*

\bar{F}_{ji}^k , \bar{F}_{bi}^c and \bar{F}_{qi}^r are given by (3.14) if in these formulae the connection coefficients are overlined.

THEOREM 5.3. *The torsion tensor $T(X, Y)$ of d -connection $\bar{\nabla}$ has the form $T(X, Y) = \bar{T}^k \delta_k + \bar{T}^c \delta_c + \bar{T}^r \delta_r$, where*

$$\begin{aligned}\bar{T}^k &= \bar{T}_{ji}^k Y^j X^i + \bar{T}_{jb}^k Y^j X^b + \bar{T}_{bj}^k Y^b X^j + \bar{T}_{jq}^k Y^j X^q + \bar{T}_{qj}^k Y^q X^j, \\ \bar{T}^c &= \bar{T}_{ji}^c Y^j X^i + \bar{T}_{ib}^c Y^j X^b + \bar{T}_{bj}^c Y^b X^j + \bar{T}_{ab}^c Y^a X^b + \bar{T}_{bq}^c Y^b X^q + \bar{T}_{qb}^c Y^q X^b, \\ \bar{T}^r &= \bar{T}_{ji}^r Y^j X^i + \bar{T}_{ba}^r Y^b X^a + \bar{T}_{jq}^r Y^j X^q + \bar{T}_{qj}^r Y^q X^j \\ &\quad + \bar{T}_{bq}^r Y^b X^q + \bar{T}_{qb}^r Y^q X^b + \bar{T}_{pq}^r Y^p X^q\end{aligned}$$

and

$$(5.2) \quad \begin{aligned}\bar{T}_{ji}^k &= \bar{F}_{ji}^k - \bar{F}_{ij}^k, & \bar{T}_{jb}^k &= -\bar{T}_{bj}^k = \bar{C}_{jb}^k, \\ \bar{T}_{jq}^k &= -\bar{T}_{qj}^k = \bar{L}_{jq}^k, & \bar{T}_{ji}^c &= -K_{ij}^c, \\ \bar{T}_{bj}^c &= -\bar{T}_{jb}^c = \bar{F}_{bj}^c - K_{jb}^c, & \bar{T}_{ab}^c &= \bar{C}_{ab}^c - \bar{C}_{ba}^c, \\ \bar{T}_{br}^c &= -\bar{T}_{rb}^c = \bar{L}_{br}^c, & \bar{T}_{qj}^r &= -\bar{T}_{jq}^r = \bar{F}_{jq}^r - K_{jq}^r, \\ \bar{T}_{qb}^r &= -\bar{T}_{bq}^r = \bar{C}_{qb}^r, & \bar{T}_{pq}^r &= \bar{L}_{pq}^r - \bar{L}_{qp}^r, \\ \bar{T}_{ji}^r &= -K_{ij}^r, & \bar{T}_{ba}^r &= -K_{ab}^r.\end{aligned}$$

It is easy to see, that all components of the torsion tensor T which appear in (5.2) with respect to (1.1) transform as tensors.

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