

# Recurrence formulas for Macdonald polynomials of type $A$

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**Abstract** We consider products of two Macdonald polynomials of type  $A$ , indexed by dominant weights which are respectively a multiple of the first fundamental weight and a weight having zero component on the  $k$ th fundamental weight. We give the explicit decomposition of any Macdonald polynomial of type  $A$  in terms of this basis.

**Keywords** Macdonald polynomials · Pieri formula · Multidimensional matrix inverse

## 1 Introduction

In the 1980s, Macdonald [6–8] introduced a class of orthogonal polynomials which are Laurent polynomials in several variables and generalize the Weyl characters of compact simple Lie groups. In the simplest situation, given a root system  $R$ , these polynomials are elements of the group algebra of the weight lattice of  $R$ , indexed by the dominant weights and depending on two parameters  $(q, t)$ .

When  $R$  is of type  $A_n$ , these Macdonald polynomials are in bijective correspondence with the symmetric functions  $\mathcal{P}_\lambda(q, t)$  indexed by partitions, introduced by Macdonald some years before [4, 5]. In fact, they correspond to  $\mathcal{P}_\lambda(q, t)(x_1, \dots, x_{n+1})$ , for a partition  $\lambda = (\lambda_1, \dots, \lambda_n)$  of length  $n$ , with the  $n + 1$  variables  $(x_1, \dots, x_{n+1})$  linked by the condition  $x_1 \cdots x_{n+1} = 1$ .

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The purpose of this article is to extend the result of [3], given for the symmetric functions  $\mathcal{P}_\lambda(q, t)$ , to the framework of the root system  $A_n$ .

More precisely, in [3, Theorem 4.1] we obtained a recurrence formula giving the symmetric function  $\mathcal{P}_{(\lambda_1, \dots, \lambda_n)}(q, t)$  as a sum

$$\mathcal{P}_{(\lambda_1, \dots, \lambda_n)} = \sum_{\theta \in \mathbb{N}^{n-1}} C_{\theta_1, \dots, \theta_{n-1}} \mathcal{P}_{(\lambda_1 + \theta_1, \dots, \lambda_{n-1} + \theta_{n-1})} \mathcal{P}_{\lambda_n - |\theta|} \tag{1.1}$$

with  $|\theta| = \sum_{i=1}^{n-1} \theta_i$  and  $\mathbb{N}$  the set of nonnegative integers. This formula was obtained by inverting the ‘‘Pieri formula,’’ which conversely expresses the product  $\mathcal{P}_{(\lambda_1, \dots, \lambda_{n-1})} \mathcal{P}_{\lambda_n}$  as a sum

$$\mathcal{P}_{(\lambda_1, \dots, \lambda_{n-1})} \mathcal{P}_{\lambda_n} = \sum_{\theta \in \mathbb{N}^{n-1}} c_{\theta_1, \dots, \theta_{n-1}} \mathcal{P}_{(\lambda_1 + \theta_1, \dots, \lambda_{n-1} + \theta_{n-1}, \lambda_n - |\theta|)}.$$

Both expansions are identities between symmetric functions, valid for any number of variables.

These identities may also be written in terms of Macdonald polynomials of type  $A_n$ . For this purpose, let  $\{\omega_i, 1 \leq i \leq n\}$  be the  $n$  fundamental weights of the root system  $A_n$ . Let  $P_\lambda$  denote the Macdonald polynomial associated with the dominant weight  $\lambda = \sum_{i=1}^n \lambda_i \omega_i$ . The recurrence formula (1.1), written for  $n + 1$  variables  $(x_1, \dots, x_{n+1})$  linked by  $x_1 \cdots x_{n+1} = 1$ , yields

$$P_\lambda = \sum_{\theta \in \mathbb{N}^{n-1}} C_{\theta_1, \dots, \theta_{n-1}} P_{(\lambda_n - |\theta|)\omega_1} P_\mu \tag{1.2}$$

with  $\mu = \sum_{i=1}^{n-2} (\lambda_i + \theta_i - \theta_{i+1})\omega_i + (\lambda_{n-1} + \lambda_n + \theta_{n-1})\omega_{n-1}$ . This alternative formulation is obvious and does not bring anything new.

However the method of [3], when applied in the  $A_n$  root system framework, allows us to get a much stronger result. Indeed, let  $k$  be a fixed integer with  $1 \leq k \leq n$ . In this paper we shall write the Macdonald polynomial  $P_\lambda$  in terms of products  $P_{r\omega_1} P_\mu$  with  $\mu = \sum_{i=1}^n \mu_i \omega_i$  and  $\mu_k = 0$ . There are  $n$  such recurrence formulas, (1.2) being the particular case  $k = n$  of the latter.

This paper is organized as follows. In Sect. 2 we introduce our notation for the root system  $A_n$  and recall general facts about the corresponding Macdonald polynomials. Their Pieri formula, which involves a specific infinite-multidimensional matrix, is studied in Sect. 3, starting from the one given by Macdonald for the symmetric functions  $\mathcal{P}_\lambda(q, t)$  [5, p. 340]. In Sect. 4 we invert the Pieri matrix by applying a particular multidimensional matrix inverse, given separately in the Appendix. This matrix inverse is equivalent to one previously obtained in [3, Sect. 2] by using operator methods. As a result of inverting the Pieri formula, we obtain recurrence formulas for  $A_n$  Macdonald polynomials. Finally, in Sect. 5 we detail the examples of the  $A_2$  and  $A_3$  cases and compare them to earlier results.

## 2 Macdonald polynomials of type A

The standard references for Macdonald polynomials associated with root systems are [6–8].

Let us consider the space  $\mathbb{R}^{n+1}$  endowed with the usual scalar product and the quotient space  $V = \mathbb{R}^{n+1}/\mathbb{R}(1, \dots, 1)$ , where  $\mathbb{R}(1, \dots, 1)$  is the subspace spanned by the vector  $(1, \dots, 1)$ . Let  $\varepsilon_1, \dots, \varepsilon_{n+1}$  denote the images in  $V$  of the coordinate vectors of  $\mathbb{R}^{n+1}$ , linked by  $\sum_{i=1}^{n+1} \varepsilon_i = 0$ .

The root system of type  $A_n$  is formed by the vectors  $\{\varepsilon_i - \varepsilon_j, i \neq j\}$ . The positive roots are  $\{\varepsilon_i - \varepsilon_j, i < j\}$ , and the simple roots are  $\varepsilon_i - \varepsilon_{i+1}$  for  $1 \leq i \leq n$ . The Weyl group is the symmetric group  $W = S_{n+1}$  acting by permutation of the coordinates.

The weight lattice  $P$  is formed by integral linear combinations of the fundamental weights  $\{\omega_i, 1 \leq i \leq n\}$  defined by  $\omega_i = \varepsilon_1 + \dots + \varepsilon_i$ . Let  $\omega_i = 0$  for  $i = 0, n + 1$ . We denote by  $P^+$  the set of dominant weights  $\lambda = \sum_{i=1}^n \lambda_i \omega_i$ , which are nonnegative integral linear combinations of the fundamental weights.

There is the following correspondence between dominant weights and partitions. Given a dominant weight, if we write it as

$$\lambda = \sum_{i=1}^n \lambda_i \omega_i = \sum_{i=1}^{n+1} \mu_i \varepsilon_i,$$

the sequence  $\mu = (\mu_1, \dots, \mu_{n+1})$  is a partition with length  $\leq n + 1$ . We have

$$\lambda_i = \mu_i - \mu_{i+1} \quad \text{and} \quad \mu_i = \mu_{n+1} + \sum_{j=i}^n \lambda_j.$$

Thus  $\mu$  is defined up to  $\mu_{n+1}$ , and two partitions  $\mu, \nu$  correspond to the same weight  $\lambda$  if and only if  $\mu_1 - \nu_1 = \dots = \mu_{n+1} - \nu_{n+1}$ . We denote by  $C_\lambda$  the family of partitions thus defined.

Let  $A$  denote the group algebra over  $\mathbb{R}$  of the free Abelian group  $P$ . For each  $\lambda \in P$ , let  $e^\lambda$  denote the corresponding element of  $A$ , subject to the multiplication rule  $e^\lambda e^\mu = e^{\lambda+\mu}$ . The set  $\{e^\lambda, \lambda \in P\}$  forms an  $\mathbb{R}$ -basis of  $A$ .

The Weyl group  $W = S_{n+1}$  acts on  $P$  and on  $A$ . Let  $W\lambda$  denote the orbit of  $\lambda \in P$  and  $A^W$  the subspace of  $W$ -invariants in  $A$ . There are two important bases of  $A^W$ , both indexed by dominant weights. The first one is given by the orbit-sums

$$m_\lambda = \sum_{\mu \in W\lambda} e^\mu.$$

The second one is provided by the Weyl characters

$$\chi_\lambda = \delta^{-1} \sum_{w \in W} \det(w) e^{w(\lambda+\rho)}$$

with  $\rho = \sum_{i=1}^n (n - i + 1)\varepsilon_i$  and  $\delta = \sum_{w \in W} \det(w) e^{w(\rho)}$ . The Macdonald polynomials  $\{P_\lambda, \lambda \in P^+\}$  form another basis defined as the eigenvectors of a specific self-adjoint operator (which we do not describe here).

For  $1 \leq i \leq n + 1$ , define  $x_i = e^{\varepsilon_i}$ , so that the variables  $x_i$  are linked by  $x_1 \cdots x_{n+1} = 1$ . Then  $\delta$  is the Vandermonde determinant  $\prod_{i < j} (x_i - x_j)$ . There is a correspondence between  $A^W$  and the symmetric polynomials restricted to  $n + 1$  variables  $x = (x_1, \dots, x_{n+1})$  linked by the previous condition.

In terms of bases this correspondence may be described as follows. Let  $\lambda$  be any dominant weight, and let  $x_1 \cdots x_{n+1} = 1$ . All monomial symmetric functions  $m_\mu(x_1, \dots, x_{n+1})$  with  $\mu \in \mathcal{C}_\lambda$  are equal, and their common value is the orbit-sum  $m_\lambda$ . Similarly, the Weyl character  $\chi_\lambda$  is the common value of the Schur functions  $s_\mu(x_1, \dots, x_{n+1})$ ,  $\mu \in \mathcal{C}_\lambda$ , whereas the Macdonald polynomial  $P_\lambda$  is the common value of the symmetric polynomials  $\mathcal{P}_\mu(q, t)(x_1, \dots, x_{n+1})$  with  $\mu \in \mathcal{C}_\lambda$  and  $\mathcal{P}_\mu(q, t)$  the symmetric function studied in Chap. 6 of [5].

Given a positive integer  $r$  and a dominant weight  $\lambda$ , the ‘‘Pieri formula’’ expands the product

$$P_{r\omega_1} P_\lambda = \sum_{\rho} c_\rho P_{\lambda+\rho}$$

in terms of Macdonald polynomials, where the range of  $\rho$  and the values of the coefficients  $c_\rho$  are to be determined.

Let  $Q$  denote the root lattice spanned by the simple roots. For any vector  $\tau$ , define

$$\Sigma(\tau) = C(\tau) \cap (\tau + Q)$$

with  $C(\tau)$  the convex hull of the Weyl group orbit of  $\tau$ . Since the orbit of  $\omega_1 = \varepsilon_1$  is the set  $\{\varepsilon_i = \omega_i - \omega_{i-1}, 1 \leq i \leq n + 1\}$ , it is clear that  $\Sigma(r\omega_1)$  is formed by vectors

$$\sum_{i=1}^{n+1} \theta_i (\omega_i - \omega_{i-1}) = \sum_{i=1}^n (\theta_i - \theta_{i+1}) \omega_i$$

with  $\theta = (\theta_1, \dots, \theta_{n+1}) \in \mathbb{N}^{n+1}$  and  $|\theta| = \sum_{i=1}^{n+1} \theta_i = r$ .

By general results [8, (5.3.8), p. 104], it is known that the sum on the right-hand side of the Pieri formula is restricted to vectors  $\rho$  such that  $\rho \in \Sigma(r\omega_1)$  and  $\lambda + \rho \in P^+$ . In the next section we shall give a direct proof of this result and make the value of the coefficient  $c_\rho$  explicit.

### 3 Pieri formula

Let  $0 < q < 1$ . For any integer  $r$ , the classical  $q$ -shifted factorial  $(u; q)_r$  is defined by

$$(u; q)_\infty = \prod_{j \geq 0} (1 - uq^j), \quad (u; q)_r = (u; q)_\infty / (uq^r; q)_\infty.$$

Let  $u = (u_1, \dots, u_m)$  be  $m$  indeterminates, and  $\theta = (\theta_1, \dots, \theta_m) \in \mathbb{N}^m$ . For clarity of display, throughout this paper, any time such a pair  $(u, \theta)$  is given, we shall implicitly assume  $m$  auxiliary variables  $v = (v_1, \dots, v_m)$  to be defined by  $v_i = q^{\theta_i} u_i$ .

Macdonald polynomials of type  $A_n$  satisfy the following Pieri formula.

**Theorem 3.1** *Let  $\lambda = \sum_{i=1}^n \lambda_i \omega_i$  be a dominant weight, and  $r \in \mathbb{N}$ . For any  $1 \leq i \leq n + 1$ , define*

$$u_i = q^{\sum_{j=i}^n \lambda_j} t^{-i}$$

and for  $\theta \in \mathbb{N}^{n+1}$ ,

$$d_\theta(u_1, \dots, u_{n+1}; r) = \frac{(q; q)_r}{(t; q)_r} \prod_{j=1}^{n+1} \frac{(t; q)_{\theta_j}}{(q; q)_{\theta_j}} \prod_{1 \leq i < j \leq n+1} \frac{(tv_i/v_j; q)_{\theta_j}}{(qv_i/v_j; q)_{\theta_j}} \frac{(qu_i/tv_j; q)_{\theta_j}}{(u_i/v_j; q)_{\theta_j}}.$$

We have

$$P_{r\omega_1} P_\lambda = \sum_{\substack{\theta \in \mathbb{N}^{n+1} \\ |\theta|=r}} d_\theta(u_1, \dots, u_{n+1}; r) P_{\lambda+\rho}$$

with  $\rho = \sum_{i=1}^n (\theta_i - \theta_{i+1}) \omega_i$ .

*Proof* In a first step, we write the Pieri formula for arbitrary  $\mathcal{P}_\mu(q, t)$  with  $\mu = (\mu_1, \dots, \mu_n)$  being a partition having length  $\leq n$ . We start from [5, p. 340, (6.24)(i)] and [5, p. 342, Example 2(a)]. Replacing  $g_r$  by  $(t; q)_r / (q; q)_r \mathcal{P}_{(r)}$ , we have

$$\mathcal{P}_{(r)} \mathcal{P}_\mu = \sum_{\kappa \supset \mu} \varphi_{\kappa/\mu} \mathcal{P}_\kappa,$$

where the skew-diagram  $\kappa - \mu$  is a horizontal  $r$ -strip, i.e., has at most one node in each column. The Pieri coefficient  $\varphi_{\kappa/\mu}$  is given by

$$\begin{aligned} \frac{(t; q)_r}{(q; q)_r} \varphi_{\kappa/\mu} &= \prod_{1 \leq i \leq j \leq l(\kappa)} \frac{f(q^{\kappa_i - \kappa_j} t^{j-i})}{f(q^{\kappa_i - \mu_j} t^{j-i})} \frac{f(q^{\mu_i - \mu_{j+1}} t^{j-i})}{f(q^{\mu_i - \kappa_{j+1}} t^{j-i})} \\ &= \prod_{1 \leq i \leq j \leq l(\kappa)} \frac{w_{\kappa_j - \mu_j} (q^{\kappa_i - \kappa_j} t^{j-i})}{w_{\kappa_{j+1} - \mu_{j+1}} (q^{\mu_i - \kappa_{j+1}} t^{j-i})} \end{aligned}$$

with  $f(u) = (tu; q)_\infty / (qu; q)_\infty$  and  $w_s(u) = (tu; q)_s / (qu; q)_s$ .

Since  $\kappa - \mu$  is a horizontal strip, the length  $l(\kappa)$  of  $\kappa$  is at most equal to  $n + 1$ , so we can write  $\kappa = (\mu_1 + \theta_1, \dots, \mu_n + \theta_n, \theta_{n+1})$  with  $|\theta| = r$ . Then

$$\begin{aligned} \frac{(t; q)_r}{(q; q)_r} \varphi_{\kappa/\mu} &= \prod_{1 \leq i \leq j \leq l(\kappa)} w_{\theta_j} (q^{\kappa_i - \kappa_j} t^{j-i}) \prod_{1 \leq i < j \leq l(\kappa)+1} (w_{\theta_j} (q^{\mu_i - \kappa_j} t^{j-i-1}))^{-1} \\ &= \prod_{j=1}^{n+1} \frac{(t; q)_{\theta_j}}{(q; q)_{\theta_j}} \prod_{1 \leq i < j \leq n+1} \frac{(tv_i/v_j; q)_{\theta_j}}{(qv_i/v_j; q)_{\theta_j}} \frac{(qu_i/tv_j; q)_{\theta_j}}{(u_i/v_j; q)_{\theta_j}}, \end{aligned}$$

where for  $1 \leq i \leq n + 1$ , we set  $u_i = q^{\mu_i} t^{-i}$  and  $v_i = q^{\kappa_i} t^{-i} = q^{\theta_i} u_i$ .

In a second step we translate this result in terms of  $A_n$  Macdonald polynomials. Given the dominant weight  $\lambda$ , we choose  $\mu = (\mu_1, \dots, \mu_{n+1})$  to be the unique

element of  $\mathcal{C}_\lambda$  such that  $\mu_{n+1} = 0$ , i.e., with length  $\leq n$ . For  $1 \leq i \leq n$ , we have  $\mu_i = \sum_{j=i}^n \lambda_j$ . As for the partition  $\kappa$  (with length  $\leq n + 1$ ), it belongs to  $\mathcal{C}_\sigma$  with  $\sigma = \sum_{k=1}^n (\kappa_k - \kappa_{k+1})\omega_k = \sum_{k=1}^n (\lambda_k + \theta_k - \theta_{k+1})\omega_k$ . Hence the statement.  $\square$

*Remark* On the right-hand side of the Pieri formula, the condition  $\lambda + \rho \in P^+$  is necessarily satisfied as soon as  $d_\theta(u_1, \dots, u_{n+1}; r) \neq 0$ . Using the correspondence between dominant weights and partitions, this may be verified on the Pieri formula

$$\mathcal{P}_{(r)}\mathcal{P}_\mu = \sum_{\kappa=(\mu_1+\theta_1, \dots, \mu_n+\theta_n, \theta_{n+1})} \varphi_{\kappa/\mu} \mathcal{P}_\kappa.$$

We only have to show that  $\varphi_{\kappa/\mu}$  necessarily vanishes when the multiinteger  $\kappa$  is not a partition. But then there is an index  $i$  such that  $\kappa_i < \kappa_{i+1}$ , so that the factor  $(qu_i/tv_{i+1}; q)_{\theta_{i+1}}$  in  $\varphi_{\kappa/\mu}$  writes out as

$$(1 - q^{1+\mu_i-\kappa_{i+1}}) \dots (1 - q^{\mu_i-\mu_{i+1}}).$$

Due to  $\kappa_i < \kappa_{i+1}$ , this product would be  $\neq 0$  only if  $\mu_i < \mu_{i+1}$ , which is impossible since  $\mu$  is a partition.

From now on, we fix some integer  $1 \leq k \leq n$ . Substituting  $r - |\theta|$  for  $\theta_k$ , the Pieri formula may be written in the more explicit form

$$P_{r\omega_1} P_\lambda = \sum_{\substack{\theta=(\theta_1, \dots, \theta_{k-1}, 0, \theta_{k+1}, \dots, \theta_{n+1}) \in \mathbb{N}^n \\ |\theta| \leq r}} \hat{d}_\theta(u_1, \dots, u_{n+1}; r) P_{\lambda+\rho}$$

with

$$\rho = \sum_{\substack{1 \leq i \leq n \\ i \neq k-1, k}} (\theta_i - \theta_{i+1})\omega_i + \theta_{k-1}\omega_{k-1} + (r - |\theta|)(\omega_k - \omega_{k-1}) - \theta_{k+1}\omega_k$$

and

$$\begin{aligned} \hat{d}_\theta(u_1, \dots, u_{n+1}; r) &= \frac{(q; q)_r}{(t; q)_r} \frac{(t; q)_{r-|\theta|}}{(q; q)_{r-|\theta|}} \prod_{\substack{j=1 \\ j \neq k}}^{n+1} \frac{(t; q)_{\theta_j}}{(q; q)_{\theta_j}} \\ &\times \prod_{\substack{1 \leq i < j \leq n+1 \\ j \neq k}} \frac{(tv_i/v_j; q)_{\theta_j}}{(qv_i/v_j; q)_{\theta_j}} \frac{(qu_i/tv_j; q)_{\theta_j}}{(u_i/v_j; q)_{\theta_j}} \\ &\times \prod_{i=1}^{k-1} \frac{(tv_i/v_k; q)_{r-|\theta|}}{(qv_i/v_k; q)_{r-|\theta|}} \frac{(qu_i/tv_k; q)_{r-|\theta|}}{(u_i/v_k; q)_{r-|\theta|}}. \end{aligned}$$

Here  $u_i, v_i$  ( $1 \leq i \leq n + 1$ ) are as in Theorem 3.1, except  $v_k = q^{r-|\theta|}u_k$ . The sum is restricted to  $|\theta| \leq r$  since  $1/(q; q)_s = 0$  for  $s < 0$ .

In a second step, we concentrate on the situation  $\lambda_k = 0$ . Then each term on the right-hand side vanishes unless  $\theta_{k+1} = 0$ . Indeed, if  $\lambda_k = 0$ , one has  $u_k = tu_{k+1}$  and  $v_{k+1} = q^{\theta_{k+1}}u_{k+1}$ . Hence, for  $i = k$  and  $j = k + 1$ , the factor  $(qu_i/tv_j; q)_{\theta_j}$  evaluates as

$$(qu_k/tv_{k+1}; q)_{\theta_{k+1}} = (q^{1-\theta_{k+1}}; q)_{\theta_{k+1}} = \delta_{\theta_{k+1},0}.$$

Therefore, if  $\lambda_k = 0$ , the Pieri formula can be written as

$$P_{r\omega_1} P_\lambda = \sum_{\substack{\theta=(\theta_1,\dots,\theta_{k-1},0,0,\theta_{k+2},\dots,\theta_{n+1}) \in \mathbb{N}^{n-1} \\ |\theta| \leq r}} \tilde{d}_\theta(u_1, \dots, u_{k-1}, u_k, u_{k+2}, \dots, u_{n+1}; k, r) P_{\lambda+\rho}$$

with

$$\rho = \sum_{\substack{1 \leq i \leq n \\ i \neq k-1, k, k+1}} (\theta_i - \theta_{i+1})\omega_i + \theta_{k-1}\omega_{k-1} + (r - |\theta|)(\omega_k - \omega_{k-1}) - \theta_{k+2}\omega_{k+1}$$

and

$$\begin{aligned} &\tilde{d}_\theta(u_1, \dots, u_{k-1}, u_k, u_{k+2}, \dots, u_{n+1}; k, r) \\ &= \frac{(q; q)_r (t; q)_{r-|\theta|}}{(t; q)_r (q; q)_{r-|\theta|}} \prod_{\substack{i=1 \\ i \neq k, k+1}}^{n+1} \frac{(t; q)_{\theta_i}}{(q; q)_{\theta_i}} \prod_{\substack{1 \leq i < j \leq n+1 \\ i \neq k, k+1 \\ j \neq k, k+1}} \frac{(tv_i/v_j; q)_{\theta_j} (qu_i/tv_j; q)_{\theta_j}}{(qv_i/v_j; q)_{\theta_j} (u_i/v_j; q)_{\theta_j}} \\ &\times \prod_{i=1}^{k-1} \frac{(tv_i/v_k; q)_{r-|\theta|} (qu_i/tv_k; q)_{r-|\theta|}}{(qv_i/v_k; q)_{r-|\theta|} (u_i/v_k; q)_{r-|\theta|}} \prod_{j=k+2}^{n+1} \frac{(tv_k/v_j; q)_{\theta_j} (qu_k/t^2v_j; q)_{\theta_j}}{(qv_k/v_j; q)_{\theta_j} (u_k/tv_j; q)_{\theta_j}}. \end{aligned}$$

Here the notation is the same as before, including  $v_k = q^{r-|\theta|}u_k$ . For  $j \geq k + 2$ , we have used

$$\begin{aligned} &\frac{(tv_k/v_j; q)_{\theta_j} (qu_k/tv_j; q)_{\theta_j} (tv_{k+1}/v_j; q)_{\theta_j} (qu_{k+1}/tv_j; q)_{\theta_j}}{(qv_k/v_j; q)_{\theta_j} (u_k/v_j; q)_{\theta_j} (qv_{k+1}/v_j; q)_{\theta_j} (u_{k+1}/v_j; q)_{\theta_j}} \\ &= \frac{(tv_k/v_j; q)_{\theta_j} (qu_k/t^2v_j; q)_{\theta_j}}{(qv_k/v_j; q)_{\theta_j} (u_k/tv_j; q)_{\theta_j}}, \end{aligned}$$

which is a direct consequence of  $v_{k+1} = u_{k+1} = u_k/t$ .

In a third step, we perform some relabeling in order to remove the two 0's appearing in  $\theta$ . For that purpose, for  $n$  indeterminates  $(u_0, u_1, \dots, u_{n-1})$  and  $\theta = (\theta_1, \dots, \theta_{n-1}) \in \mathbb{N}^{n-1}$ , we define

$$\begin{aligned} &D_\theta(u_0, u_1, \dots, u_{n-1}; k, r) \\ &= (q/t)^{|\theta|} \frac{(t^2u_0; q)_{|\theta|}}{(qtu_0; q)_{|\theta|}} \prod_{i=1}^{n-1} \frac{(t; q)_{\theta_i} (q^{|\theta|+1}u_i; q)_{\theta_i}}{(q; q)_{\theta_i} (q^{|\theta|}tu_i; q)_{\theta_i}} \end{aligned}$$

$$\begin{aligned} & \times \prod_{1 \leq i < j \leq n-1} \frac{(tv_i/v_j; q)_{\theta_j}}{(qv_i/v_j; q)_{\theta_j}} \frac{(qu_i/tv_j; q)_{\theta_j}}{(u_i/v_j; q)_{\theta_j}} \\ & \times \prod_{i=1}^{k-1} \frac{(u_i/u_0; q)_{\theta_i}}{(qu_i/tu_0; q)_{\theta_i}} \frac{(qu_i/tu_0; q)_{\theta_i-r+|\theta|}}{(u_i/u_0; q)_{\theta_i-r+|\theta|}} \frac{(u_i/tu_0; q)_{\theta_i-r+|\theta|}}{(qu_i/t^2u_0; q)_{\theta_i-r+|\theta|}} \\ & \times \prod_{i=k}^{n-1} \frac{(tu_i/u_0; q)_{\theta_i}}{(qu_i/u_0; q)_{\theta_i}}. \end{aligned}$$

**Lemma** *If we write*

$$w_i = \begin{cases} q^{-r}t^{-2}, & i = 0, \\ q^{-r}u_i/tu_k, & 1 \leq i \leq k-1, \\ q^{-r}u_{i+2}/tu_k, & k \leq i \leq n-1, \end{cases}$$

*we have*

$$\begin{aligned} & D_\theta(w_0, w_1, \dots, w_{n-1}; k, r) \\ & = \tilde{d}_{(\theta_1, \dots, \theta_{k-1}, 0, 0, \theta_k, \dots, \theta_{n-1})}(u_1, \dots, u_{k-1}, u_k, u_{k+2}, \dots, u_{n+1}; k, r). \end{aligned}$$

*Proof* Merely by substitution, and using  $v_k = q^{r-|\theta|}u_k$ , we only have to prove

$$\begin{aligned} & (q/t)^{|\theta|} \frac{(q^{-r}; q)_{|\theta|}}{(q^{1-r}/t; q)_{|\theta|}} \prod_{j=k+2}^{n+1} \frac{(q^{|\theta|-r+1}u_j/tu_k; q)_{\theta_j}}{(q^{|\theta|-r}u_j/u_k; q)_{\theta_j}} \frac{(t^2u_j/u_k; q)_{\theta_j}}{(qtu_j/u_k; q)_{\theta_j}} \\ & \times \prod_{i=1}^{k-1} \frac{(q^{|\theta|-r+1}u_i/tu_k; q)_{\theta_i}}{(q^{|\theta|-r}u_i/u_k; q)_{\theta_i}} \frac{(tu_i/u_k; q)_{\theta_i}}{(qu_i/u_k; q)_{\theta_i}} \\ & \times \prod_{i=1}^{k-1} \frac{(qu_i/u_k; q)_{\theta_i-r+|\theta|}}{(tu_i/u_k; q)_{\theta_i-r+|\theta|}} \frac{(u_i/u_k; q)_{\theta_i-r+|\theta|}}{(qu_i/tu_k; q)_{\theta_i-r+|\theta|}} \\ & = \frac{(q; q)_r}{(t; q)_r} \frac{(t; q)_{r-|\theta|}}{(q; q)_{r-|\theta|}} \prod_{i=1}^{k-1} \frac{(tv_i/q^{r-|\theta|}u_k; q)_{r-|\theta|}}{(qv_i/q^{r-|\theta|}u_k; q)_{r-|\theta|}} \frac{(qu_i/tq^{r-|\theta|}u_k; q)_{r-|\theta|}}{(u_i/q^{r-|\theta|}u_k; q)_{r-|\theta|}} \\ & \times \prod_{j=k+2}^{n+1} \frac{(tq^{r-|\theta|}u_k/v_j; q)_{\theta_j}}{(q^{r-|\theta|+1}u_k/v_j; q)_{\theta_j}} \frac{(qu_k/t^2v_j; q)_{\theta_j}}{(u_k/tv_j; q)_{\theta_j}}. \end{aligned}$$

We have obviously

$$\frac{(q^{|\theta|-r+1}u_i/tu_k; q)_{\theta_i}}{(q^{|\theta|-r}u_i/u_k; q)_{\theta_i}} \frac{(u_i/u_k; q)_{\theta_i-r+|\theta|}}{(qu_i/tu_k; q)_{\theta_i-r+|\theta|}} = \frac{(qu_i/tq^{r-|\theta|}u_k; q)_{r-|\theta|}}{(u_i/q^{r-|\theta|}u_k; q)_{r-|\theta|}}.$$

Using the identities

$$\frac{(aq^{-n}; q)_n}{(bq^{-n}; q)_n} = \frac{(q/a; q)_n}{(q/b; q)_n} (a/b)^n,$$



$$\frac{(a; q)_n (b; q)_{n-k}}{(b; q)_n (a; q)_{n-k}} = \frac{(q^{1-n}/a; q)_k}{(q^{1-n}/b; q)_k} (a/b)^k,$$

we get

$$\begin{aligned} \frac{(tu_i/u_k; q)_{\theta_i} (qu_i/u_k; q)_{\theta_i-r+|\theta|}}{(qu_i/u_k; q)_{\theta_i} (tu_i/u_k; q)_{\theta_i-r+|\theta|}} &= \frac{(q^{1-\theta_i}u_k/tu_i; q)_{r-|\theta|}}{(q^{-\theta_i}u_k/u_i; q)_{r-|\theta|}} (t/q)^{r-|\theta|} \\ &= \frac{(tv_i/q^{r-|\theta|}u_k; q)_{r-|\theta|}}{(qv_i/q^{r-|\theta|}u_k; q)_{r-|\theta|}}. \end{aligned}$$

Similarly, we obtain

$$\begin{aligned} (t/q)^{\theta_j} \frac{(q^{|\theta|-r+1}u_j/tu_k; q)_{\theta_j}}{(q^{|\theta|-r}u_j/u_k; q)_{\theta_j}} &= \frac{(tq^{r-|\theta|}u_k/v_j; q)_{\theta_j}}{(q^{r-|\theta|+1}u_k/v_j; q)_{\theta_j}}, \\ (q/t)^{\theta_j} \frac{(t^2u_j/u_k; q)_{\theta_j}}{(qtu_j/u_k; q)_{\theta_j}} &= \frac{(qu_k/t^2v_j; q)_{\theta_j}}{(u_k/tv_j; q)_{\theta_j}}. \end{aligned} \quad \square$$

Finally, we have proved the following Pieri formula.

**Theorem 3.2** *Let  $\lambda = \sum_{i=1}^n \lambda_i \omega_i$  be a dominant weight, and  $r \in \mathbb{N}$ . Assume  $\lambda_k = 0$  for some fixed  $1 \leq k \leq n$ . Define*

$$u_i = \begin{cases} q^{-r}t^{-2}, & i = 0, \\ q^{-r+\sum_{j=i}^{k-1} \lambda_j}t^{k-i-1}, & 1 \leq i \leq k-1, \\ q^{-r-\sum_{j=k+1}^{i+1} \lambda_j}t^{k-i-3}, & k \leq i \leq n-1. \end{cases}$$

We have

$$P_{r\omega_1} P_\lambda = \sum_{\substack{\theta=(\theta_1, \dots, \theta_{n-1}) \in \mathbb{N}^{n-1} \\ |\theta| \leq r}} D_\theta(u_0, u_1, \dots, u_{n-1}; k, r) P_{\lambda+\rho}$$

with

$$\begin{aligned} \rho &= \sum_{i=1}^{k-2} (\theta_i - \theta_{i+1})\omega_i + \theta_{k-1}\omega_{k-1} + (r - |\theta|)(\omega_k - \omega_{k-1}) - \theta_k\omega_{k+1} \\ &+ \sum_{i=k+2}^n (\theta_{i-2} - \theta_{i-1})\omega_i. \end{aligned}$$

*Remark* For  $k = 1, 2$  (resp.  $k = n, n - 1$ ), the first (resp. the last) sum in the above expression of  $\rho$  must be understood as zero. This convention will be kept in the next sections.

### 4 A recurrence formula

Given two multiintegers  $\beta = (\beta_1, \dots, \beta_{n-1})$  and  $\kappa = (\kappa_1, \dots, \kappa_{n-1}) \in \mathbb{Z}^{n-1}$ , we write  $\beta \geq \kappa$  for  $\beta_i \geq \kappa_i$  ( $1 \leq i \leq n - 1$ ). We say that an infinite  $(n - 1)$ -dimensional matrix  $F = (f_{\beta\kappa})_{\beta, \kappa \in \mathbb{Z}^{n-1}}$  is lower-triangular if  $f_{\beta\kappa} = 0$  unless  $\beta \geq \kappa$ . When all  $f_{\kappa\kappa} \neq 0$ , there exists a unique lower-triangular matrix  $G = (g_{\kappa\gamma})_{\kappa, \gamma \in \mathbb{Z}^{n-1}}$  such that

$$\sum_{\beta \geq \kappa \geq \gamma} f_{\beta\kappa} g_{\kappa\gamma} = \delta_{\beta\gamma}$$

for all  $\beta, \gamma \in \mathbb{Z}^{n-1}$ , where  $\delta_{\beta\gamma}$  is the usual Kronecker symbol. We refer to  $F$  and  $G$  as mutually inverse.

Such a pair of infinite multidimensional inverse matrices is given in the [Appendix](#), as a corollary of [3, Theorem 2.7] (and, in fact, equivalent to the latter). This result is essential for our purpose.

Given  $n$  indeterminates  $(u_0, u_1, \dots, u_{n-1})$ ,  $\theta = (\theta_1, \dots, \theta_{n-1}) \in \mathbb{N}^{n-1}$ , and  $k, r \in \mathbb{N}$  with  $1 \leq k \leq n$ , we define

$$\begin{aligned} & C_{\theta_1, \dots, \theta_{n-1}}(u_0, u_1, \dots, u_{n-1}; k, r) \\ &= q^{|\theta|} \frac{(t^2 u_0; q)_{|\theta|}}{(qtu_0; q)_{|\theta|}} \prod_{i=1}^{n-1} \frac{(q/t; q)_{\theta_i}}{(q; q)_{\theta_i}} \frac{(qu_i; q)_{\theta_i}}{(qtu_i; q)_{\theta_i}} \\ &\quad \times \prod_{1 \leq i < j \leq n-1} \frac{(qv_i/tv_j; q)_{\theta_j}}{(qv_i/v_j; q)_{\theta_j}} \frac{(tu_i/v_j; q)_{\theta_j}}{(u_i/v_j; q)_{\theta_j}} \\ &\quad \times \prod_{i=1}^{k-1} \frac{(u_i/tu_0; q)_{\theta_i}}{(qu_i/t^2u_0; q)_{\theta_i}} \frac{(qtu_0/u_i; q)_r}{(t^2u_0/u_i; q)_r} \frac{(tu_0/u_i; q)_r}{(qu_0/u_i; q)_r} \prod_{i=k}^{n-1} \frac{(tu_i/u_0; q)_{\theta_i}}{(qu_i/u_0; q)_{\theta_i}} \\ &\quad \times \frac{1}{\Delta(v)} \det_{1 \leq i, j \leq n-1} \left[ v_i^{n-j-1} \left( 1 - t^{j-1} \frac{1-tv_i}{1-v_i} \prod_{s=1}^{n-1} \frac{v_i - u_s}{v_i - tv_s} \right) \right] \end{aligned}$$

with  $\Delta(v)$  the Vandermonde determinant  $\prod_{1 \leq i < j \leq n-1} (v_i - v_j)$ . Here is our main result.

**Theorem 4.1** *Let  $\lambda = \sum_{i=1}^n \lambda_i \omega_i$  be a dominant weight. Assume  $\lambda_k = 0$  for some fixed  $1 \leq k \leq n$ . For any positive integer  $r \leq \lambda_{k-1}$ , the weight*

$$\lambda^{(r)} = \lambda + r(\omega_k - \omega_{k-1}) = \lambda + r\epsilon_k$$

is dominant. Define

$$u_i = \begin{cases} q^{-r} t^{-2}, & i = 0, \\ q^{-r + \sum_{j=i}^{k-1} \lambda_j} t^{k-i-1}, & 1 \leq i \leq k-1, \\ q^{-r - \sum_{j=k+1}^{i+1} \lambda_j} t^{k-i-3}, & k \leq i \leq n-1. \end{cases}$$

We have

$$P_{\lambda(r)} = \sum_{\substack{\theta=(\theta_1, \dots, \theta_{n-1}) \in \mathbb{N}^{n-1} \\ |\theta| \leq r}} C_{\theta}(u_0, u_1, \dots, u_{n-1}; k, r) P_{(r-|\theta)\omega_1} P_{\lambda+\rho}$$

with

$$\rho = \sum_{i=1}^{k-2} (\theta_i - \theta_{i+1})\omega_i + \theta_{k-1}\omega_{k-1} - \theta_k\omega_{k+1} + \sum_{i=k+2}^n (\theta_{i-2} - \theta_{i-1})\omega_i.$$

*Remark* The weight  $\lambda + \rho$  has no component on  $\omega_k$ . Further, similarly as in Theorem 3.1 (see the remark following the proof of that theorem), the condition  $\lambda + \rho \in P^+$  is necessarily satisfied in Theorem 4.2 as soon as  $C_{\theta}(u_0, u_1, \dots, u_{n-1}; k, r) \neq 0$ . We omit the details which involve a tedious case-by-case analysis.

*Proof* We make use of the multidimensional matrix inverse given in the Appendix. Let  $\beta = (\beta_1, \dots, \beta_{n-1}), \kappa = (\kappa_1, \dots, \kappa_{n-1}), \gamma = (\gamma_1, \dots, \gamma_{n-1}) \in \mathbb{Z}^{n-1}$ . If we define

$$f_{\beta\kappa} = C_{\beta_1-\kappa_1, \dots, \beta_{n-1}-\kappa_{n-1}}(q^{|\kappa|}u_0, q^{\kappa_1+|\kappa|}u_1, \dots, q^{\kappa_{n-1}+|\kappa|}u_{n-1}; k, r - |\kappa|),$$

$$g_{\kappa\gamma} = D_{\kappa_1-\gamma_1, \dots, \kappa_{n-1}-\gamma_{n-1}}(q^{|\gamma|}u_0, q^{\gamma_1+|\gamma|}u_1, \dots, q^{\gamma_{n-1}+|\gamma|}u_{n-1}; k, r - |\gamma|),$$

by this result, the infinite lower-triangular multidimensional matrices  $(f_{\beta\kappa})_{\beta, \kappa \in \mathbb{Z}^{n-1}}$  and  $(g_{\kappa\gamma})_{\kappa, \gamma \in \mathbb{Z}^{n-1}}$  are mutually inverse.

Now let us replace  $\lambda_i$  in Theorem 3.2 by  $\lambda_i + \gamma_i - \gamma_{i+1}$  for  $1 \leq i \leq k - 2$ ,  $\lambda_{k-1}$  by  $\lambda_{k-1} + \gamma_{k-1}$ ,  $\lambda_{k+1}$  by  $\lambda_{k+1} - \gamma_k$ ,  $\lambda_i$  by  $\lambda_i + \gamma_{i-2} - \gamma_{i-1}$  for  $k + 2 \leq i \leq n$ , and  $r$  by  $r - |\gamma|$ . Then  $u_0$  is replaced by  $q^{|\gamma|}u_0$ , and  $u_i$  by  $q^{\gamma_i+|\gamma|}u_i$  for  $1 \leq i \leq n - 1$ . In explicit terms, we are considering the identity

$$P_{(r-|\gamma)\omega_1} P_{\lambda+\tilde{\gamma}}$$

$$= \sum_{\substack{\theta=(\theta_1, \dots, \theta_{n-1}) \in \mathbb{N}^{n-1} \\ |\theta| \leq r}} D_{\theta}(q^{|\gamma|}u_0, q^{\gamma_1+|\gamma|}u_1, \dots, q^{\gamma_{n-1}+|\gamma|}u_{n-1}; k, r - |\gamma|) P_{\lambda+\tilde{\gamma}+\rho}$$

with

$$u_i = \begin{cases} q^{-r}t^{-2}, & i = 0, \\ q^{-r+\sum_{j=i}^{k-1} \lambda_j}t^{k-i-1}, & 1 \leq i \leq k - 1, \\ q^{-r-\sum_{j=k+1}^{i+1} \lambda_j}t^{k-i-3}, & k \leq i \leq n - 1, \end{cases}$$

and

$$\rho = \sum_{i=1}^{k-2} (\theta_i - \theta_{i+1})\omega_i + \theta_{k-1}\omega_{k-1} + (r - |\theta|)(\omega_k - \omega_{k-1}) - \theta_k\omega_{k+1}$$

$$+ \sum_{i=k+2}^n (\theta_{i-2} - \theta_{i-1})\omega_i,$$

$$\tilde{\gamma} = \sum_{i=1}^{k-2} (\gamma_i - \gamma_{i+1})\omega_i + \gamma_{k-1}\omega_{k-1} - \gamma_k\omega_{k+1} + \sum_{i=k+2}^n (\gamma_{i-2} - \gamma_{i-1})\omega_i.$$

After substituting the summation indices  $\theta_i \mapsto \kappa_i - \gamma_i$  for  $1 \leq i \leq n - 1$ , we obtain exactly

$$\sum_{\kappa \in \mathbb{Z}^{n-1}} g_{\kappa\gamma} y_{\kappa} = w_{\gamma} \quad (\gamma \in \mathbb{Z}^{n-1})$$

with

$$y_{\kappa} = P_{\lambda+\tilde{\kappa}}, \quad w_{\gamma} = P_{(r-|\gamma|)\omega_1} P_{\lambda+\tilde{\gamma}}$$

and

$$\begin{aligned} \tilde{\kappa} &= \sum_{i=1}^{k-2} (\kappa_i - \kappa_{i+1})\omega_i + \kappa_{k-1}\omega_{k-1} + (r - |\kappa|)(\omega_k - \omega_{k-1}) - \kappa_k\omega_{k+1} \\ &+ \sum_{i=k+2}^n (\kappa_{i-2} - \kappa_{i-1})\omega_i. \end{aligned}$$

This immediately yields the inverse relation

$$\sum_{\beta \in \mathbb{Z}^{n-1}} f_{\beta\kappa} w_{\beta} = y_{\kappa} \quad (\kappa \in \mathbb{Z}^{n-1}).$$

We conclude by setting  $\kappa_i = 0$  for all  $1 \leq i \leq n - 1$ . □

Finally, by the substitutions  $r \rightarrow \lambda_k$  and  $\lambda_{k-1} \rightarrow \lambda_{k-1} + \lambda_k$ , we obtain the following very remarkable expansion.

**Theorem 4.2** *Let  $\lambda = \sum_{i=1}^n \lambda_i \omega_i$  be a dominant weight, and  $k \in \mathbb{N}$  fixed with  $1 \leq k \leq n$ . Define*

$$u_i = \begin{cases} q^{-\lambda_k} t^{-2}, & i = 0, \\ q^{\sum_{j=i}^{k-1} \lambda_j} t^{k-i-1}, & 1 \leq i \leq k - 1, \\ q^{-\sum_{j=k}^{i+1} \lambda_j} t^{k-i-3}, & k \leq i \leq n - 1, \end{cases}$$

and  $\mu = \lambda - \lambda_k (\omega_k - \omega_{k-1}) = \lambda - \lambda_k \varepsilon_k$ . We have

$$P_{\lambda} = \sum_{\substack{\theta=(\theta_1, \dots, \theta_{n-1}) \in \mathbb{N}^{n-1} \\ |\theta| \leq \lambda_k}} C_{\theta}(u_0, u_1, \dots, u_{n-1}; k, \lambda_k) P_{(\lambda_k - |\theta|)\omega_1} P_{\mu+\rho}$$

with

$$\rho = \sum_{i=1}^{k-2} (\theta_i - \theta_{i+1})\omega_i + \theta_{k-1}\omega_{k-1} - \theta_k\omega_{k+1} + \sum_{i=k+2}^n (\theta_{i-2} - \theta_{i-1})\omega_i.$$

*Remark* Observe that the weights  $\mu$  and  $\mu + \rho$  have no component on  $\omega_k$ .

The special case  $k = n$  is worth writing out explicitly.

**Corollary** Let  $\lambda = \sum_{i=1}^n \lambda_i \omega_i$  be a dominant weight. Define  $u_0 = q^{-\lambda_n} t^{-2}$  and  $u_i = q^{\sum_{l=i}^{n-1} \lambda_l} t^{n-i-1}$  ( $1 \leq i \leq n - 1$ ). We have

$$P_\lambda = \sum_{\substack{\theta=(\theta_1, \dots, \theta_{n-1}) \in \mathbb{N}^{n-1} \\ |\theta| \leq \lambda_n}} C_\theta(u_0, u_1, \dots, u_{n-1}; n, \lambda_n) P_{(\lambda_n - |\theta|)\omega_1} P_\mu$$

with  $\mu = \sum_{i=1}^{n-2} (\lambda_i + \theta_i - \theta_{i+1})\omega_i + (\lambda_{n-1} + \lambda_n + \theta_{n-1})\omega_{n-1}$ .

The reader may check that this is exactly Theorem 4.1 of [3] (with  $n \mapsto n - 1$ ), written for  $x_1 \cdots x_{n+1} = 1$ , up to the normalization  $Q_\lambda = b_\lambda P_\lambda$  with

$$b_\lambda = \prod_{1 \leq i \leq j \leq n} \frac{(q^{\sum_{l=i}^{j-1} \lambda_l} t^{j-i+1}; q)_{\lambda_j}}{(q^{1+\sum_{l=i}^{j-1} \lambda_l} t^{j-i}; q)_{\lambda_j}} = \prod_{1 \leq i \leq j \leq n} \frac{(tu_i/u_j; q)_{\lambda_j}}{(qu_i/u_j; q)_{\lambda_j}},$$

where we set  $u_n = 1/t$ .

### 5 Examples

In this section we write out the formulas in Theorem 4.2 explicitly for  $n = 2, 3$ .

#### 5.1 The root system $A_2$

For  $k = 2$ , we have  $u_0 = q^{-\lambda_2} / t^2$ ,  $u_1 = q^{\lambda_1}$ , and

$$C_\theta(u_0, u_1; 2, r) = q^\theta \frac{(t^2 u_0; q)_\theta}{(q t u_0; q)_\theta} \frac{(q/t; q)_\theta}{(q; q)_\theta} \frac{(q u_1; q)_\theta}{(q t u_1; q)_\theta} \frac{(u_1 / t u_0; q)_\theta}{(q u_1 / t^2 u_0; q)_\theta} \\ \times \frac{(q t u_0 / u_1; q)_r}{(t^2 u_0 / u_1; q)_r} \frac{(t u_0 / u_1; q)_r}{(q u_0 / u_1; q)_r} \left( 1 - \frac{1 - t v_1}{1 - v_1} \frac{v_1 - u_1}{v_1 - t u_1} \right).$$

After some simplifications, we obtain

$$P_{\lambda_1 \omega_1 + \lambda_2 \omega_2} = \sum_{\theta \in \mathbb{N}} C_\theta^{(2)}(\lambda) P_{(\lambda_2 - \theta)\omega_1} P_{(\lambda_1 + \lambda_2 + \theta)\omega_1}$$

with

$$C_\theta^{(2)}(\lambda) = C_\theta(u_0, u_1; 2, \lambda_2) \\ = t^\theta \frac{(q^{\lambda_2 - \theta + 1}; q)_\theta}{(t q^{\lambda_2 - \theta}; q)_\theta} \frac{(1/t; q)_\theta}{(q; q)_\theta} \frac{(q^{\lambda_1 + 1}; q)_\theta}{(t q^{\lambda_1 + 1}; q)_\theta}$$

$$\times \frac{(tq^{\lambda_1}; q)_{\lambda_2+\theta}}{(q^{\lambda_1+1}; q)_{\lambda_2+\theta}} \frac{(tq^{\lambda_1+1}; q)_{\lambda_2}}{(t^2q^{\lambda_1}; q)_{\lambda_2}} \frac{1 - q^{\lambda_1+2\theta}}{1 - q^{\lambda_1+\theta}}.$$

This result may be compared with the Jing–Józefiak classical result [1], more precisely, with its restriction to three variables  $(x_1, x_2, x_3)$  subject to  $x_1x_2x_3 = 1$ . Namely, given a partition  $(\mu_1, \mu_2)$ , the Macdonald symmetric function  $\mathcal{P}_{(\mu_1, \mu_2)}(q, t)$  is given by

$$\mathcal{P}_{(\mu_1, \mu_2)} = \sum_{\theta \in \mathbb{N}} C_\theta(\mu) \mathcal{P}_{(\mu_2-\theta)} \mathcal{P}_{(\mu_1+\theta)}$$

with

$$C_\theta(\mu) = \frac{(tq^{\mu_1-\mu_2+1}; q)_{\mu_2}}{(t^2q^{\mu_1-\mu_2}; q)_{\mu_2}} \frac{(q^{\mu_2-\theta+1}; q)_\theta}{(tq^{\mu_2-\theta}; q)_\theta} \frac{(tq^{\mu_1-\mu_2}; q)_{\mu_2+\theta}}{(q^{\mu_1-\mu_2+1}; q)_{\mu_2+\theta}} \\ \times t^\theta \frac{(1/t; q)_\theta}{(q; q)_\theta} \frac{(q^{\mu_1-\mu_2+1}; q)_\theta}{(tq^{\mu_1-\mu_2+1}; q)_\theta} \frac{1 - q^{\mu_1-\mu_2+2\theta}}{1 - q^{\mu_1-\mu_2+\theta}}.$$

Our formula is equivalent to the main result of [1] by the correspondence  $\lambda_1 = \mu_1 - \mu_2, \lambda_2 = \mu_2$  between dominant weights and partitions, recalled in Sect. 2.

For  $k = 1$ , we have  $u_0 = q^{-\lambda_1}/t^2, u_1 = q^{-\lambda_1-\lambda_2}/t^3$ , and

$$C_\theta(u_0, u_1; 1, r) = q^\theta \frac{(t^2u_0; q)_\theta}{(qtu_0; q)_\theta} \frac{(q/t; q)_\theta}{(q; q)_\theta} \frac{(qu_1; q)_\theta}{(qtu_1; q)_\theta} \frac{(tu_1/u_0; q)_\theta}{(qu_1/u_0; q)_\theta} \\ \times \left( 1 - \frac{1 - tv_1}{1 - v_1} \frac{v_1 - u_1}{v_1 - tu_1} \right).$$

After some simplifications, we obtain

$$P_{\lambda_1\omega_1+\lambda_2\omega_2} = \sum_{\theta \in \mathbb{N}} C_\theta^{(1)}(\lambda) P_{(\lambda_1-\theta)\omega_1} P_{(\lambda_2-\theta)\omega_2}$$

with

$$C_\theta^{(1)}(\lambda) = C_\theta(u_0, u_1; 1, \lambda_1) \\ = t^\theta \frac{(1/t; q)_\theta}{(q; q)_\theta} \frac{(q^{\lambda_1}; 1/q)_\theta}{(tq^{\lambda_1-1}; 1/q)_\theta} \frac{(q^{\lambda_2}; 1/q)_\theta}{(tq^{\lambda_2-1}; 1/q)_\theta} \frac{(t^3q^{\lambda_1+\lambda_2-1}; 1/q)_\theta}{(t^2q^{\lambda_1+\lambda_2-1}; 1/q)_\theta} \\ \times \frac{1 - t^3q^{\lambda_1+\lambda_2-2\theta}}{1 - t^3q^{\lambda_1+\lambda_2-\theta}}.$$

We thus recover exactly the result of Perelomov, Ragoucy, and Zaugg [9, Theorem 1(a)].

### 5.2 The root system $A_3$

For  $k = 1, 2, 3$ , our formulas in Theorem 4.2 write respectively as

$$P_{\lambda_1\omega_1+\lambda_2\omega_2+\lambda_3\omega_3} = \sum_{(i,j) \in \mathbb{N}^2} C_{ij}^{(1)}(\lambda) P_{(\lambda_1-i-j)\omega_1} P_{(\lambda_2-i)\omega_2+(\lambda_3+i-j)\omega_3}$$

$$\begin{aligned}
 &= \sum_{(i,j) \in \mathbb{N}^2} C_{ij}^{(2)}(\lambda) P_{(\lambda_2-i-j)\omega_1} P_{(\lambda_1+\lambda_2+i)\omega_1+(\lambda_3-j)\omega_3} \\
 &= \sum_{(i,j) \in \mathbb{N}^2} C_{ij}^{(3)}(\lambda) P_{(\lambda_3-i-j)\omega_1} P_{(\lambda_1+i-j)\omega_1+(\lambda_2+\lambda_3+j)\omega_2}.
 \end{aligned}$$

In order to make these expansions explicit, we need to evaluate the determinant of the 2 by 2 matrix  $A$  given by

$$A_{kl} = v_k^{2-l} \left( 1 - t^{l-1} \frac{1 - tv_k}{1 - v_k} \frac{v_k - u_1}{v_k - tu_1} \frac{v_k - u_2}{v_k - tu_2} \right)$$

with  $v_1 = q^i u_1, v_2 = q^j u_2$ .

More precisely, we need to compute the quotient of this determinant by the Vandermonde determinant  $v_1 - v_2 = q^i u_1 - q^j u_2$ . There is no evidence that this quotient may be written in canonical form. Inspired by the explicit result of [2, Theorem 1] (see below), we write this quotient of determinants as

$$\begin{aligned}
 \frac{\det A}{q^i u_1 - q^j u_2} &= \frac{(t - 1)^2}{(t - q^i)(t - q^j)} \\
 &\times \left( \frac{1 - q^{2i} u_1}{1 - q^i u_1} \frac{1 - q^{2j} u_2}{1 - q^j u_2} \left( 1 + t^{-1} \frac{1 - q^i}{1 - q^i u_1 / tu_2} \frac{1 - q^j}{1 - q^j u_2 / tu_1} \right) \right. \\
 &\left. - (q^i u_1 + q^j u_2) \frac{1 - q^i}{1 - q^i u_1} \frac{1 - q^j}{1 - q^j u_2} \frac{1 - q^i / t}{1 - q^i u_1 / tu_2} \frac{1 - q^j / t}{1 - q^j u_2 / tu_1} \right).
 \end{aligned}$$

The above identity (which is not trivial) may be easily verified by using any formal calculus software.

Next, for  $(i, j) \in \mathbb{N}^2$ , we define

$$\begin{aligned}
 &\nabla_{ij}(u_0, u_1, u_2) \\
 &= q^{i+j} \frac{(t^2 u_0; q)_{i+j}}{(qtu_0; q)_{i+j}} \frac{(1/t; q)_i}{(q; q)_i} \frac{(u_1; q)_i}{(qtu_1; q)_i} \frac{(1/t; q)_j}{(q; q)_j} \\
 &\times \frac{(u_2; q)_j}{(qtu_2; q)_j} \frac{(q^{i-j+1} u_1 / tu_2; q)_j}{(q^{i-j+1} u_1 / u_2; q)_j} \frac{(tq^{-j} u_1 / u_2; q)_j}{(q^{-j} u_1 / u_2; q)_j} \\
 &\times \left( \frac{1 - q^{2i} u_1}{1 - u_1} \frac{1 - q^{2j} u_2}{1 - u_2} \left( 1 + t^{-1} \frac{1 - q^i}{1 - q^i u_1 / tu_2} \frac{1 - q^j}{1 - q^j u_2 / tu_1} \right) \right. \\
 &\left. - (q^i u_1 + q^j u_2) \frac{1 - q^i}{1 - u_1} \frac{1 - q^j}{1 - u_2} \frac{1 - q^i / t}{1 - q^i u_1 / tu_2} \frac{1 - q^j / t}{1 - q^j u_2 / tu_1} \right).
 \end{aligned}$$

It is readily verified that we have

$$\frac{C_{ij}(u_0, u_1, u_2; 1, r)}{\nabla_{ij}(u_0, u_1, u_2)} = \frac{(tu_1/u_0; q)_i (tu_2/u_0; q)_j}{(qu_1/u_0; q)_i (qu_2/u_0; q)_j},$$

$$\frac{C_{ij}(u_0, u_1, u_2; 2, r)}{\nabla_{ij}(u_0, u_1, u_2)} = \frac{(u_1/tu_0; q)_i}{(qu_1/t^2u_0; q)_i} \frac{(qtu_0/u_1; q)_r}{(t^2u_0/u_1; q)_r} \frac{(tu_0/u_1; q)_r}{(qu_0/u_1; q)_r} \frac{(tu_2/u_0; q)_j}{(qu_2/u_0; q)_j},$$

$$\frac{C_{ij}(u_0, u_1, u_2; 3, r)}{\nabla_{ij}(u_0, u_1, u_2)} = \frac{(u_1/tu_0; q)_i}{(qu_1/t^2u_0; q)_i} \frac{(qtu_0/u_1; q)_r}{(t^2u_0/u_1; q)_r} \frac{(tu_0/u_1; q)_r}{(qu_0/u_1; q)_r}$$

$$\times \frac{(u_2/tu_0; q)_j}{(qu_2/t^2u_0; q)_j} \frac{(qtu_0/u_2; q)_r}{(t^2u_0/u_2; q)_r} \frac{(tu_0/u_2; q)_r}{(qu_0/u_2; q)_r}.$$

Now, by Theorem 4.2 the respective recurrence coefficients are determined to be

$$C_{ij}^{(1)}(\lambda) = C_{ij}(q^{-\lambda_1}/t^2, q^{-\lambda_1-\lambda_2}/t^3, q^{-\lambda_1-\lambda_2-\lambda_3}/t^4; 1, \lambda_1),$$

$$C_{ij}^{(2)}(\lambda) = C_{ij}(q^{-\lambda_2}/t^2, q^{\lambda_1}, q^{-\lambda_2-\lambda_3}/t^3; 2, \lambda_2),$$

$$C_{ij}^{(3)}(\lambda) = C_{ij}(q^{-\lambda_3}/t^2, q^{\lambda_1+\lambda_2}t, q^{\lambda_2}; 3, \lambda_3).$$

The cases  $k = 1, 2$  are new. For  $k = 3$ , we recover the first author’s earlier result in [2, Theorem 1], more precisely, the restriction of this result to four variables  $(x_1, x_2, x_3, x_4)$  subject to  $x_1x_2x_3x_4 = 1$ . Namely, given a partition  $(\mu_1, \mu_2, \mu_3)$  and  $u = q^{\mu_1-\mu_2}$ ,  $v = q^{\mu_2-\mu_3}$ , the Macdonald symmetric function  $\mathcal{P}_{(\mu_1, \mu_2, \mu_3)}(q, t)$  is given by

$$\mathcal{P}_{(\mu_1, \mu_2, \mu_3)} = \sum_{(i, j) \in \mathbb{N}^2} C_{ij}(\mu) \mathcal{P}_{(\mu_3-i-j)} \mathcal{P}_{(\mu_1+i, \mu_2+j)}$$

with

$$C_{ij}(\mu) = t^{i+j} \frac{(1/t; q)_i}{(q; q)_i} \frac{(1/t; q)_j}{(q; q)_j} \frac{(tuv; q)_i}{(qt^2uv; q)_i} \frac{(v; q)_j}{(qtv; q)_j} \frac{(q^{-j}t^2u; q)_i}{(q^{-j}tu; q)_i} \frac{(qu; q)_i}{(qtu; q)_i}$$

$$\times \frac{(t; q)_{\mu_1-\mu_2+i-j}}{(q; q)_{\mu_1-\mu_2+i-j}} \frac{(t; q)_{\mu_2+j}}{(q; q)_{\mu_2+j}} \frac{(t; q)_{\mu_3-i-j}}{(q; q)_{\mu_3-i-j}} \frac{(q; q)_{\mu_1-\mu_2}}{(t; q)_{\mu_1-\mu_2}}$$

$$\times \frac{(q; q)_{\mu_2-\mu_3}}{(t; q)_{\mu_2-\mu_3}} \frac{(q; q)_{\mu_3}}{(t; q)_{\mu_3}} \frac{(q^{i-j}t^2u; q)_{\mu_2+j}}{(q^{i-j+1}tu; q)_{\mu_2+j}} \frac{(qtu; q)_{\mu_2-\mu_3}}{(t^2u; q)_{\mu_2-\mu_3}}$$

$$\times \frac{(qt^2uv; q)_{\mu_3}}{(t^3uv; q)_{\mu_3}} \frac{(qtv; q)_{\mu_3}}{(t^2v; q)_{\mu_3}} \frac{1 - q^{2i}tuv}{1 - tuv} \frac{1 - q^{2j}v}{1 - v}$$

$$\times \left( 1 + u \frac{1 - q^i}{1 - q^i u} \frac{1 - q^{-j}}{1 - q^{-j} t^2 u} \left( t - v(q^i tu + q^j) \frac{t - q^i}{1 - q^{2i} tuv} \frac{t - q^j}{1 - q^{2j} v} \right) \right).$$

The reader may check our formula is indeed equivalent to [2, Theorem 1] by using the correspondence  $\lambda_1 = \mu_1 - \mu_2, \lambda_2 = \mu_2 - \mu_3, \lambda_3 = \mu_3$  between dominant weights and partitions.

### 6 Final remark

The Macdonald polynomial  $P_\lambda, \lambda = \sum_{i=1}^n \lambda_i \omega_i$ , is in bijective correspondence with the symmetric function  $\mathcal{P}_\mu(x_1, \dots, x_{n+1})$  with  $\mu = (\mu_1, \dots, \mu_n), \mu_i = \sum_{j=i}^n \lambda_j,$



subject to the condition  $x_1 \cdots x_{n+1} = 1$ . Therefore the  $n$  recurrence relations that we have obtained for  $P_\lambda$  may be expressed in terms of  $\mathcal{P}_\mu(x_1, \dots, x_{n+1})$ , subject to  $x_1 \cdots x_{n+1} = 1$ .

One may wonder whether this restriction can be removed. Equivalently, given some fixed integer  $1 \leq k \leq n$ , is it possible to expand the symmetric function  $\mathcal{P}_\mu$  in terms of products  $\mathcal{P}_{(r)}\mathcal{P}_\rho$  for partitions  $\rho = (\rho_1, \dots, \rho_n)$  satisfying  $\rho_k = \rho_{k+1}$ ?

Such a development has been obtained in [3] for  $k = n$ , in which case  $\rho_n = \rho_{n+1} = 0$ . However, this method cannot be used for other values of  $k$ .

Actually, the Pieri expansion of  $\mathcal{P}_{(r)}\mathcal{P}_\rho$  involves symmetric functions  $\mathcal{P}_\sigma$  with  $\sigma - \rho$  a horizontal  $r$ -strip. Hence some of these partitions  $\sigma$  have length  $l(\sigma) = n + 1$ . The only exception occurs for  $k = n$  since in that case  $\rho_n = 0$  entails  $l(\sigma) \leq n$ .

Therefore, except for  $k = n$ , the Pieri multiplication does not conserve the space generated by  $\{\mathcal{P}_\kappa, l(\kappa) \leq n\}$ , and it is not possible to define a Pieri matrix to invert.

This difficulty does not arise in the  $A_n$  framework. Then the Pieri matrix can be defined, because the condition  $x_1 \cdots x_{n+1} = 1$  and the property [5, (4.17), p. 325]

$$\mathcal{P}_{(\sigma_1, \dots, \sigma_{n+1})}(x_1, \dots, x_{n+1}) = (x_1 \cdots x_{n+1})^{\sigma_{n+1}} \mathcal{P}_{(\sigma_1 - \sigma_{n+1}, \dots, \sigma_n - \sigma_{n+1}, 0)}(x_1, \dots, x_{n+1})$$

allow us to deal with partitions of length  $n + 1$ .

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**Appendix: A multidimensional matrix inverse**

The following result (equivalent to one previously given in [3]) is crucial to obtain the recursion formula in Sect. 4.

**Lemma** *Let  $t, u_0, u_1, \dots, u_n$  be indeterminates, and  $r, k \in \mathbb{N}$  with  $1 \leq k \leq n + 1$ . Define*

$$\begin{aligned} f_{\beta\kappa} &= q^{|\beta| - |\kappa|} \frac{(t^2u_0; q)_{|\beta|}}{(qtu_0; q)_{|\beta|}} \frac{(qtu_0; q)_{|\kappa|}}{(t^2u_0; q)_{|\kappa|}} \prod_{i=1}^n \frac{(q/t; q)_{\beta_i - \kappa_i}}{(q; q)_{\beta_i - \kappa_i}} \frac{(q^{\kappa_i + |\kappa| + 1}u_i; q)_{\beta_i - \kappa_i}}{(q^{\kappa_i + |\kappa| + 1}tu_i; q)_{\beta_i - \kappa_i}} \\ &\times \prod_{i=1}^{k-1} \frac{(u_i/tu_0; q)_{\beta_i}}{(qu_i/t^2u_0; q)_{\beta_i}} \frac{(qu_i/tu_0; q)_{\kappa_i}}{(u_i/u_0; q)_{\kappa_i}} \\ &\times \prod_{i=1}^{k-1} \frac{(u_i/u_0; q)_{|\kappa| - r + \kappa_i}}{(qu_i/tu_0; q)_{|\kappa| - r + \kappa_i}} \frac{(qu_i/t^2u_0; q)_{|\kappa| - r + \kappa_i}}{(u_i/tu_0; q)_{|\kappa| - r + \kappa_i}} \\ &\times \prod_{i=k}^n \frac{(tu_i/u_0; q)_{\beta_i}}{(qu_i/u_0; q)_{\beta_i}} \frac{(qu_i/u_0; q)_{\kappa_i}}{(tu_i/u_0; q)_{\kappa_i}} \\ &\times \prod_{1 \leq i < j \leq n} \frac{(q^{\beta_i - \beta_j + 1}u_i/tu_j; q)_{\beta_j - \kappa_j}}{(q^{\beta_i - \beta_j + 1}u_i/u_j; q)_{\beta_j - \kappa_j}} \frac{(q^{\kappa_i - \beta_j}tu_i/u_j; q)_{\beta_j - \kappa_j}}{(q^{\kappa_i - \beta_j}u_i/u_j; q)_{\beta_j - \kappa_j}} \frac{1}{q^{\beta_i}u_i - q^{\beta_j}u_j} \end{aligned}$$

$$\times \det_{1 \leq i, j \leq n} \left[ (q^{\beta_i} u_i)^{n-j} \left( 1 - t^{j-1} \frac{(1 - q^{\beta_i + |\kappa|} t u_i)}{(1 - q^{\beta_i + |\kappa|} u_i)} \prod_{s=1}^n \frac{(q^{\beta_i} u_i - q^{\kappa_s} u_s)}{(q^{\beta_i} u_i - q^{\kappa_s} t u_s)} \right) \right]$$

and

$$\begin{aligned} g_{\kappa\gamma} &= \left(\frac{q}{t}\right)^{|\kappa| - |\gamma|} \frac{(t^2 u_0; q)_{|\kappa|}}{(qtu_0; q)_{|\kappa|}} \frac{(qtu_0; q)_{|\gamma|}}{(t^2 u_0; q)_{|\gamma|}} \prod_{i=1}^n \frac{(t; q)_{\kappa_i - \gamma_i}}{(q; q)_{\kappa_i - \gamma_i}} \frac{(q^{\gamma_i + |\kappa| + 1} u_i; q)_{\kappa_i - \gamma_i}}{(q^{\gamma_i + |\kappa|} t u_i; q)_{\kappa_i - \gamma_i}} \\ &\times \prod_{i=1}^{k-1} \frac{(u_i/u_0; q)_{\kappa_i}}{(qu_i/tu_0; q)_{\kappa_i}} \frac{(qu_i/t^2 u_0; q)_{\gamma_i}}{(u_i/tu_0; q)_{\gamma_i}} \\ &\times \prod_{i=1}^{k-1} \frac{(qu_i/tu_0; q)_{|\kappa| - r + \kappa_i}}{(u_i/u_0; q)_{|\kappa| - r + \kappa_i}} \frac{(u_i/tu_0; q)_{|\kappa| - r + \kappa_i}}{(qu_i/t^2 u_0; q)_{|\kappa| - r + \kappa_i}} \\ &\times \prod_{i=k}^n \frac{(tu_i/u_0; q)_{\kappa_i}}{(qu_i/u_0; q)_{\kappa_i}} \frac{(qu_i/u_0; q)_{\gamma_i}}{(tu_i/u_0; q)_{\gamma_i}} \\ &\times \prod_{1 \leq i < j \leq n} \frac{(q^{\kappa_i - \kappa_j} t u_i/u_j; q)_{\kappa_j - \gamma_j}}{(q^{\kappa_i - \kappa_j + 1} u_i/u_j; q)_{\kappa_j - \gamma_j}} \frac{(q^{\gamma_i - \kappa_j + 1} u_i/tu_j; q)_{\kappa_j - \gamma_j}}{(q^{\gamma_i - \kappa_j} u_i/u_j; q)_{\kappa_j - \gamma_j}}. \end{aligned}$$

Then the infinite lower-triangular  $n$ -dimensional matrices  $(f_{\beta\kappa})_{\beta, \kappa \in \mathbb{Z}^n}$  and  $(g_{\kappa\gamma})_{\kappa, \gamma \in \mathbb{Z}^n}$  are mutually inverse.

*Proof* Given two nonzero sequences  $(\xi_\kappa)$  and  $(\zeta_\kappa)$  and a pair of matrices  $(f_{\beta\kappa})$  and  $(g_{\kappa\gamma})$  which are mutually inverse, it is easily checked (using the trivial relation  $\frac{\xi_\beta}{\xi_\gamma} \delta_{\beta\gamma} = \delta_{\beta\gamma}$ ) that the matrices  $(f_{\beta\kappa} \xi_\beta / \zeta_\kappa)$  and  $(g_{\kappa\gamma} \zeta_\kappa / \xi_\gamma)$  are mutually inverse.

We choose

$$\begin{aligned} \xi_\kappa &= \left(\frac{q}{t}\right)^{|\kappa|} \frac{(t^2 u_0; q)_{|\kappa|}}{(qtu_0; q)_{|\kappa|}} \prod_{i=1}^{k-1} \frac{(u_i/tu_0; q)_{\kappa_i}}{(qu_i/t^2 u_0; q)_{\kappa_i}} \prod_{i=k}^n \frac{(tu_i/u_0; q)_{\kappa_i}}{(qu_i/u_0; q)_{\kappa_i}} \\ &\times \prod_{1 \leq i < j \leq n} \frac{(qu_i/u_j; q)_{\kappa_i - \kappa_j}}{(tu_i/u_j; q)_{\kappa_i - \kappa_j}} \frac{(u_i/u_j; q)_{\kappa_i - \kappa_j}}{(qu_i/tu_j; q)_{\kappa_i - \kappa_j}}, \\ \zeta_\kappa &= \left(\frac{q}{t}\right)^{|\kappa|} \frac{(t^2 u_0; q)_{|\kappa|}}{(qtu_0; q)_{|\kappa|}} \prod_{i=1}^{k-1} \frac{(u_i/u_0; q)_{\kappa_i}}{(qu_i/tu_0; q)_{\kappa_i}} \\ &\times \prod_{i=1}^{k-1} \frac{(qu_i/tu_0; q)_{|\kappa| - r + \kappa_i}}{(u_i/u_0; q)_{|\kappa| - r + \kappa_i}} \frac{(u_i/tu_0; q)_{|\kappa| - r + \kappa_i}}{(qu_i/t^2 u_0; q)_{|\kappa| - r + \kappa_i}} \\ &\times \prod_{i=k}^n \frac{(tu_i/u_0; q)_{\kappa_i}}{(qu_i/u_0; q)_{\kappa_i}} \prod_{1 \leq i < j \leq n} \frac{(qu_i/u_j; q)_{\kappa_i - \kappa_j}}{(tu_i/u_j; q)_{\kappa_i - \kappa_j}} \frac{(u_i/u_j; q)_{\kappa_i - \kappa_j}}{(qu_i/tu_j; q)_{\kappa_i - \kappa_j}}, \end{aligned}$$

together with the pair of mutually inverse matrices  $(f_{\beta\kappa})$  and  $(g_{\kappa\gamma})$  as defined in [3, Theorem 2.7].

Several elementary manipulations of  $q$ -shifted factorials eventually lead to the result in the desired form. To give a sample (concentrating only on the products over  $\prod_{1 \leq i < j \leq n}$  of  $q$ -shifted factorials), we use the simplification

$$\begin{aligned} & \frac{(q^{\kappa_i - \kappa_j + 1} u_i / t u_j; q)_{\beta_i - \kappa_i} (q^{\kappa_i - \beta_j} t u_i / u_j; q)_{\beta_i - \kappa_i}}{(q^{\kappa_i - \kappa_j + 1} u_i / u_j; q)_{\beta_i - \kappa_i} (q^{\kappa_i - \beta_j} u_i / u_j; q)_{\beta_i - \kappa_i}} \\ & \quad \times \frac{(q u_i / u_j; q)_{\beta_i - \beta_j} (u_i / u_j; q)_{\beta_i - \beta_j} (t u_i / u_j; q)_{\kappa_i - \kappa_j} (q u_i / t u_j; q)_{\kappa_i - \kappa_j}}{(t u_i / u_j; q)_{\beta_i - \beta_j} (q u_i / t u_j; q)_{\beta_i - \beta_j} (q u_i / u_j; q)_{\kappa_i - \kappa_j} (u_i / u_j; q)_{\kappa_i - \kappa_j}} \\ & = \frac{(q u_i / t u_j; q)_{\beta_i - \kappa_j} (u_i / u_j; q)_{\kappa_i - \beta_j} (q u_i / u_j; q)_{\beta_i - \beta_j} (t u_i / u_j; q)_{\kappa_i - \kappa_j}}{(q u_i / u_j; q)_{\beta_i - \kappa_j} (t u_i / u_j; q)_{\kappa_i - \beta_j} (q u_i / t u_j; q)_{\beta_i - \beta_j} (u_i / u_j; q)_{\kappa_i - \kappa_j}} \\ & = \frac{(q^{\beta_i - \beta_j + 1} u_i / t u_j; q)_{\beta_j - \kappa_j} (q^{\kappa_i - \beta_j} t u_i / u_j; q)_{\beta_j - \kappa_j}}{(q^{\beta_i - \beta_j + 1} u_i / u_j; q)_{\beta_j - \kappa_j} (q^{\kappa_i - \beta_j} u_i / u_j; q)_{\beta_j - \kappa_j}} \end{aligned}$$

in the computation of  $f_{\beta\kappa}$  in the lemma. □

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