

Tenth MSU Conference on Differential Equations and Computational Simulations.
Electronic Journal of Differential Equations, Conference 23 (2016), pp. 77–86.
 ISSN: 1072-6691. URL: <http://ejde.math.txstate.edu> or <http://ejde.math.unt.edu>
<ftp://ejde.math.txstate.edu>

MULTIPLICITY OF SOLUTIONS OF RELATIVISTIC-TYPE SYSTEMS WITH PERIODIC NONLINEARITIES: A SURVEY

JEAN MAWHIN

ABSTRACT. We survey recent results on the multiplicity of T -periodic solutions of differential systems of the form

$$\left(\frac{u'}{\sqrt{1-|u'|^2}}\right)' + \nabla_u F(t, u) = e(t)$$

when $F(t, u)$ is ω_i -periodic with respect to u_i ($i = 1, \dots, N$). Several techniques of critical point theory are used.

1. MOTIVATION: THE FORCED PENDULUM AND CORRESPONDING SYSTEMS

The periodic problem for the forced pendulum equation

$$u'' + a \sin u = e(t), \quad u(0) = u(T), \quad u'(0) = u'(T) \quad (1.1)$$

has been for almost one century a source of inspiration for ordinary differential equations and nonlinear functional analysis, and a cornerstone for most nonlinear techniques (see e.g. [16, 18]). In particular its solutions are the critical points of the Lagrangian action functional

$$\mathcal{L}(u) := \int_0^T \left[\frac{u'^2}{2} + a \cos u + eu \right] dt$$

in the Sobolev space $H_T^1 = \{u \in H^1([0, T]) : u(0) = u(T)\}$.

In 1922, Hamel [12] proved that for each $e \in C([0, T])$ such that

$$\bar{e} := T^{-1} \int_0^T e(t) dt = 0,$$

there exists at least one solution of (1.1) minimizing $\mathcal{L}(u)$ over T -periodic C^1 functions. This result was rapidly forgotten and, following a renewal of interest for the problem, in 1980, due to Castro's application [6] of some minimax method to (1.1), Hamel's theorem was rediscovered independently around 1981 by Willem [27] and Dancer [10] in the more natural framework of H_T^1 . Because of the structure of the equation, if u is a solution, the same is true for $u + 2k\pi$ for all $k \in \mathbb{Z}$, so that two

2010 *Mathematics Subject Classification.* 34C15, 34C25, 58E05.

Key words and phrases. Pendulum-type equations; multiple solutions; critical point theory;

Ljusternik-Schnirelmann category.

©2016 Texas State University.

Published March 21, 2016.

solutions of (1.1) are called *geometrically distinct* if they do not differ by an integer multiple of 2π .

As shown by the special case of the unforced pendulum, Hamel's existence conclusion is not optimal and, in 1984, the following multiplicity result was proved in [20].

Theorem 1.1. *For each $e \in L^1(0, T)$ such that $\bar{e} = 0$, problem (1.1) has at least two geometrically distinct solutions.*

The second solution was obtained by a mountain pass type argument between a minimizing solution u_0 and the other one $u_0 + 2\pi$. The unforced case shows that this multiplicity result is optimal if no restriction is made upon a and T . An immediate generalization of Theorem 1.1, based upon the same arguments, holds for $a \sin u$ replaced by a Carathéodory function $f(t, u)$ such that $F(t, u) := \int_0^u f(t, s) ds$ is ω -periodic in u for a.e. fixed $t \in [0, T]$, and some $\omega > 0$.

The solutions of the N -dimensional corresponding problem

$$u'' + \nabla_u F(t, u) = e(t), \quad u(0) = u(T), \quad u'(0) = u'(T), \quad (1.2)$$

where $e \in L^1(0, T; \mathbb{R}^N)$, $F : [0, T] \times \mathbb{R}^N \rightarrow \mathbb{R}$ and $\nabla_u F : [0, T] \times \mathbb{R}^N \rightarrow \mathbb{R}^N$ are Carathéodory functions such that

$$F(t, u + \omega_j e_j) = F(t, u) \quad (j = 1, \dots, N) \quad (1.3)$$

for a.e. $t \in [0, T]$, all $u \in \mathbb{R}^N$, and some $\omega_i > 0$ ($i = 1, \dots, N$), are the critical points of the Lagrangian action functional

$$\mathcal{L}_N(u) := \int_0^T \left[\frac{|u'|^2}{2} - F(t, u) + (e|u) \right] dt$$

in the Sobolev space $H_T^1 = \{u \in H^1([0, T], \mathbb{R}^N) : u(0) = u(T)\}$. Here and in the whole paper, $(\cdot|\cdot)$ denotes the inner product in \mathbb{R}^N and $|\cdot|$ the corresponding norm. In 1984, the following result was proved in [21].

Theorem 1.2. *If F satisfies assumption (1.3), then, for each $e \in L^1(0, T; \mathbb{R}^N)$ such that $\bar{e} = 0$, problem (1.2) has at least two geometrically distinct solutions.*

Geometrically distinct solutions of (1.2) are of course solutions whose differences are not of the form $\sum_{i=1}^N k_i \omega_i$ for some $(k_1, \dots, k_N) \in \mathbb{Z}^N$. The proof of Theorem 1.2 is an easy extension of the argument of the scalar case. Such a multiplicity result is not optimal, as easily seen, and was improved around 1988 independently by Rabinowitz [24], Chang [7], and the author [17], who got the following multiplicity conclusion.

Theorem 1.3. *If F satisfies assumption (1.3), then, for each $e \in L^1(0, T; \mathbb{R}^N)$ such that $\bar{e} = 0$, problem (1.2) has at least $N + 1$ geometrically distinct solutions.*

Although they present technical differences, the three proofs of this result use the fact that $\mathcal{L}_N(u + \omega_j e^j) = \mathcal{L}_N(u)$ ($j = 1, \dots, n$) and some Ljusternik-Schnirelmann category arguments.

2. THE RELATIVISTIC FORCED PENDULUM AND CORRESPONDING SYSTEMS

In 2010, it was shown in [4] that the solutions of the 'relativistic forced pendulum equation', i.e. the solutions of the problem

$$\left(\frac{u'}{\sqrt{1 - u'^2}} \right)' + a \sin u = e(t), \quad u(0) = u(T), \quad u'(0) = u'(T), \quad (2.1)$$

namely the functions u of class C^1 on $C([0, T])$ such that $\|u'\|_\infty < 1$, $\frac{u'}{\sqrt{1-u'^2}}$ is absolutely continuous on $[0, T]$ and which verify the differential equation in (2.1) almost everywhere, and the periodic boundary conditions, can be associated to the critical points of the action defined by

$$\mathcal{R}(u) := \int_0^T [1 - \sqrt{1 - |u'|^2} + a \cos u + eu] dt$$

on the closed convex set

$$K = \{u \in W^{1,\infty}([0, T]) : u(0) = u(T), \|u'\|_\infty \leq 1\},$$

where $\|\cdot\|_\infty$ denotes the L^∞ -norm. In 2011, it was shown in [2] that those solutions could be seen as well be associated to the critical points in the sense of Szulkin [25] of the functional given on $C_T = \{u \in C([0, T]) : u(0) = u(T)\}$ by

$$\mathcal{S}(u) = \Phi(u) + \mathcal{G}(u),$$

where Φ is defined on $C([0, T])$ by

$$\Phi(u) := \begin{cases} \int_0^T [1 - \sqrt{1 - |u'|^2}] dt & \text{if } u \in W^{1,\infty}([0, T]) \\ +\infty & \text{if } u \in C([0, T]) \setminus W^{1,\infty}([0, T]) \end{cases}$$

and \mathcal{G} is defined on $C([0, T])$ by

$$\mathcal{G}(u) = \int_0^T [a \cos u + eu] dt.$$

Φ is convex, proper, lower semi-continuous, and \mathcal{G} of class C^1 , so that \mathcal{S} has the structure required by Szulkin's critical point theory [25]. When $\bar{e} = 0$, \mathcal{R} and \mathcal{S} are bounded from below, and satisfy a suitable version of Palais-Smale condition on their set of definition. Consequently, they reach there a minimum, and one can show that such minimum corresponds to a solution of (2.1) (this is less trivial than in the case of (1.1)). Hence, the following extension of Hamel's result to (2.1) follows [4, 2] : *for each $e \in L^1(0, T)$ such that $\bar{e} = 0$, problem (2.1) has at least one solution minimizing \mathcal{R} on K (or \mathcal{S} on C_T)*. Another proof of this existence result, based upon some Hamiltonian equivalent formulation (described later in a different context) and a saddle point theorem, has been given in 2012 by Manásevich and Ward [14].

Like in the classical case, such a conclusion is not optimal for (2.1) and, in 2012, Bereanu and Torres [3] have proved the following multiplicity result.

Theorem 2.1. *For each $e \in C([0, T])$ such that $\bar{e} = 0$, problem (2.1) has at least two geometrically distinct solutions.*

Their proof is modeled on the one of [20] for the classical pendulum, but technically more involved. One first shows the existence of two positive minimizers u_0 , $u_0 + 2\pi$ of \mathcal{S} on C_T , then considers a modified problem like in the method of lower and upper solutions with lower solution $\alpha = u_0$, and upper solution $\beta = u_1$. Finally one shows that Szulkin's critical points of the corresponding modified action are solutions of (2.1), and obtains the second solution of the modified action by a mountain pass argument.

In 2012, Fonda and Toader [11] and, in 2013, Marò [15] have proved Theorem 2.1 by applying a Poincaré-Birkhoff type fixed point theorem to the equivalent Hamiltonian formulation mentioned above, and Marò has obtained the supplementary

information that *one of the solutions is unstable*. An extension of Theorem 2.1 is easily obtained when $a \sin u$ is replaced by $\partial_u F(t, u)$, with $F(t, u)$ ω -periodic in u .

If $e \in L^1(0, T; \mathbb{R}^N)$, if the functions $F(t, u)$ and $\nabla_u F(t, u)$ are defined and continuous on $[0, T] \times \mathbb{R}^N$, and if assumption (1.3) holds, one can consider the periodic problem for a relativistic system

$$\left(\frac{u'}{\sqrt{1 - |u'|^2}} \right)' + \nabla_u F(t, u) = e(t), \quad u(0) = u(T), \quad u'(0) = u'(T). \quad (2.2)$$

Its concept of solution is defined in an analogous way as for (2.1). With now

$$C_T := \{u \in C([0, T], \mathbb{R}^N) : u(0) = u(T)\},$$

one defines $\mathcal{S} : C_T \rightarrow (-\infty, +\infty]$ by

$$\mathcal{S}(u) = \Phi(u) + \mathcal{G}(u), \quad (2.3)$$

where

$$\Phi(u) := \begin{cases} \int_0^T [1 - \sqrt{1 - |u'|^2}] dt & \text{if } u \in W^{1, \infty}([0, T], \mathbb{R}^N) \\ +\infty & \text{if } u \in C_T \setminus W^{1, \infty}([0, T], \mathbb{R}^N), \end{cases}$$

and

$$\mathcal{G}(u) = \int_0^T [-F(t, u) + (e|u)] dt.$$

In 2011, using the approach of [4], the following existence result was proved in [5].

Theorem 2.2. *If F satisfies assumption (1.3), then, for each $e \in L^1(0, T; \mathbb{R}^N)$ such that $\bar{e} = 0$, problem (2.2) has at least one solution.*

The extension to systems of the methods used in [3, 11, 15] for a scalar equation seeming difficult, the obtention of multiplicity results similar to Theorem 1.3 for system (2.2) has required different approaches, that we now describe.

3. A HAMILTONIAN APPROACH

The first result was given in 2012 in [19]. It is assumed that F and $\nabla_u F$ exist and are continuous, and, for simplicity, we extend them as well as e , by T -periodicity, to $\mathbb{R} \times \mathbb{R}^N$ and to \mathbb{R} respectively. Setting

$$v = \frac{u'}{\sqrt{1 - |u'|^2}}$$

in (2.2), which is equivalent to

$$u' = \frac{v}{\sqrt{1 + |v|^2}},$$

we immediately see that (2.2) is equivalent to the first order problem

$$v' = -\nabla_u F(t, u) + e(t), \quad u' = \frac{v}{\sqrt{1 + |v|^2}}, \quad u, v \text{ } T\text{-periodic.} \quad (3.1)$$

Defining $H : \mathbb{R} \times \mathbb{R}^N \times \mathbb{R}^N \rightarrow \mathbb{R}$ by

$$H(t, u, v) := \sqrt{1 + |v|^2} - 1 + F(t, u) - (e(t)|u),$$

we see that (3.1) has the Hamiltonian form

$$v' = -\nabla_u H(t, u, v), \quad u' = \nabla_v H(t, u, v), \quad u, v \text{ } T\text{-periodic.} \quad (3.2)$$

The action functional naturally associated to (3.2) is given by

$$\mathcal{H}(v, u) = \int_0^T [-(v|u') + F(t, u) - (e|u) + \sqrt{1 + |v|^2} - 1] dt.$$

If we define the Sobolev space

$$H_T^{1/2} = \{(v, u) \in H^{1/2}([0, T], \mathbb{R}^{2N}) : v \text{ and } u \text{ are } T\text{-periodic}\},$$

then a standard result (see e.g. [23]) implies that if $e \in L^s(0, T; \mathbb{R}^N)$ for some $s > 1$, then $\mathcal{H} \in C^1(H_T^{1/2}, \mathbb{R})$ and its critical points solve (3.2), or, explicitly, (3.1). Furthermore, it is easy to check that if $\bar{e} = 0$, and F satisfies condition (1.3), then

$$\mathcal{H}(v, u_1 + k_1\omega_1, \dots, u_N + k_N\omega_N) = \mathcal{H}(v, u)$$

for all $(v, u) \in H_T^{1/2}$ and all $(k_1, \dots, k_N) \in \mathbb{Z}^N$. Consequently, we can consider \mathcal{H} as defined on $\mathbb{T}^N \times E$, where

$$E = \{(v, u) \in H_T^{1/2} : \bar{u} = 0\}$$

and \mathbb{T}^N is the n -torus. It can be shown that $E = E^- \oplus E^0 \oplus E^+$ where $E^0 \simeq \mathbb{R}^N$, the linear operator associated to the quadratic form $(v, u) \mapsto \int_0^T [-(v|u')] dt$ is negative definite on E^- , positive definite on E^+ , and

$$\mathcal{H}(v, \bar{u}) \rightarrow +\infty \quad \text{for all } \bar{u} \in \mathbb{R}^N \text{ when } |v| \rightarrow \infty \text{ in } E^0.$$

Therefore \mathcal{H} satisfies the conditions of an abstract saddle point theorem for indefinite functionals proved in 1990 by Szulkin [25], and based upon the concept of relative Ljusternik-Schnirelmann category, which implies the following multiplicity result for (2.2).

Theorem 3.1. *If F satisfies assumption (1.3), then, for each $e \in L^s(0, T; \mathbb{R}^N)$ for some $s > 1$, such that $\bar{e} = 0$, problem (2.2) has at least $N + 1$ geometrically distinct solutions.*

As one can see, the proof of Theorem 3.1 is technically sophisticated, both from the critical point theory side, because \mathcal{H} is an indefinite functional, and from the topological side, because the relative category is a more involved and delicate concept than Ljusternik-Schnirelmann category. Hence the result of [19] raises the following natural questions:

- (1) Can $e \in L^s$ for some $s > 1$ be replaced by the more natural assumption $e \in L^1$?
- (2) Can the result be proved using Lagrangian action and classical category?

4. A LAGRANGIAN APPROACH

In 2013, Bereanu and Jebelean [1] proved Theorem 3.1, when F , $\nabla_u F$ and e are continuous, through an extension to convex, lower semicontinuous perturbations of a C^1 -functional on a Banach space, i.e. to functionals of Szulkin type [25], of an abstract multiplicity result for some symmetric C^1 functionals given in [22, Theorem 4.12], and motivated by Rabinowitz' approach in [24].

Let X be a Banach space with dual X^* and duality mapping $\langle \cdot, \cdot \rangle$, G a discrete additive subgroup of X such that $\text{span}(G)$ has finite dimension N , $\pi : X \rightarrow X/G$ the canonical projection. So $G \simeq \mathbb{Z}^N$, $X = \mathbb{R}^N \oplus Y$ for some closed subspace Y , $u = \bar{u} + \tilde{u}$, with $\bar{u} \in \mathbb{R}^N$, $\tilde{u} \in Y$. $A \subset X$ is G -invariant if $u + g \in A$ for all $u \in A$

and $g \in G$, $f : X \rightarrow M$ is G -invariant if $f(u + g) = f(u)$ for all $u \in X$ and $g \in G$. The following assumptions are made:

- (H1) $\mathcal{G} \in C^1(X, \mathbb{R})$ is G -invariant, \mathcal{G}' takes bounded sets into bounded sets.
- (H2) $\Psi : X \rightarrow (-\infty, +\infty]$ is G -invariant, convex, lower semicontinuous, with closed non-empty domain $D(\Psi) \supset \{u \in X : \|\tilde{u}\| \leq \rho, |\Psi(u)| \leq \rho\}$, $\Psi(0) = 0$, $\Psi(u) = \Psi(\tilde{u})$ for all $u \in X$.
- (H3) Any sequence (u_n) in X with (\bar{u}_n) bounded has a convergent subsequence.

According to [25], $u \in X$ is a *critical point* of $\mathcal{S} = \Psi + \mathcal{G}$ if

$$\langle \mathcal{G}'(u), v - u \rangle + \Psi(v) - \Psi(u) \geq 0 \quad \text{for all } v \in X.$$

Let

$$K = \{u \in X : u \text{ is a critical point}\}$$

be the *critical set* of \mathcal{S} , and let $K_c = \{u \in K : \mathcal{S}(u) = c\}$. It is easy to see that $\mathcal{S}, \mathcal{S}', K, K_c$ are G -invariant. Hence, if u is a critical point of \mathcal{S} , the same is true for $u + g$ for all $g \in G$, and the set $\{u + g : g \in G\}$ is called a *critical orbit* of \mathcal{S} .

If \mathcal{N} is an open neighborhood of K_c and $\varepsilon > 0$, we set

$$\mathcal{N}_\varepsilon = \{u \in X \setminus \mathcal{N} : |\bar{u}| \leq 2, \mathcal{S}(u) \leq c + \varepsilon\}.$$

The following equivariant deformation lemma, which combines similar results in [22, 25], is essential to prove the multiplicity result.

Lemma 4.1. *Let $c \in \mathbb{R}$ and \mathcal{N} be a G -invariant neighborhood of K_c . Then, for each $\varepsilon \in (0, 1]$, there exists $\bar{\varepsilon} \in (0, \varepsilon]$, $d_\varepsilon > 0$, $\varepsilon' \in (0, \varepsilon]$ and $\eta \in C([0, \bar{t}] \times \mathcal{N}_\varepsilon, X)$, with the following properties:*

- (i) $\eta(0, \cdot) = id_{\mathcal{N}_\varepsilon}$.
- (ii) $\eta(t, u + g) = \eta(t, u) + g$ for all $(t, u) \in [0, \bar{t}] \times \mathcal{N}_\varepsilon$ and all $g \in G$ with $u + g \in \mathcal{N}_\varepsilon$.
- (iii) $\|\eta(t, u) - u\| \leq d_\varepsilon t$ for all $(t, u) \in [0, \bar{t}] \times \mathcal{N}_\varepsilon$.
- (iv) $\mathcal{S}(\eta(t, u)) - \mathcal{S}(u) \leq d_\varepsilon t$ for all $(t, u) \in [0, \bar{t}] \times \mathcal{N}_\varepsilon$.
- (v) $\mathcal{S}(\eta(t, u)) - \mathcal{S}(u) \leq -\varepsilon' t/2$ for all $(t, u) \in [0, \bar{t}] \times (\mathcal{N}_\varepsilon \cap \mathcal{S}^{-1}([c - \varepsilon, +\infty)))$.
- (vi) if $A \subset \mathcal{N}_\varepsilon$ with $c \leq \sup_A \mathcal{S}$, then, for all $t \in [0, \bar{t}]$,

$$\sup_A \mathcal{S}(\eta(t, \cdot)) - \sup_A \mathcal{S} \leq -\varepsilon' t/2.$$

From this lemma, one can construct a deformation in the quotient space $\pi(X)$. Defining, like in [22],

$$\mathcal{A}_j = \{A \subset X : A \text{ is compact and } \text{cat}_{\pi(X)}(\pi(A)) \geq j\},$$

one can check that $\mathcal{A}_j \neq \emptyset$ for each $j = 1, \dots, N + 1$ and \mathcal{A}_j is a complete metric space for the Hausdorff distance. Furthermore, the function $\sigma : \mathcal{A}_j \rightarrow (-\infty, +\infty]$ defined by

$$\sigma(A) = \sup_{A \in \mathcal{A}_j} \mathcal{S}$$

is lower semicontinuous and bounded from below. Ekeland's variational principle and a rather standard argument of Ljusternik-Schirelmann type give the following multiplicity result.

Proposition 4.2. *Under assumptions (H1)–(H3), the functional $\mathcal{S} = \Psi + \mathcal{G}$ has at least $N + 1$ critical orbits.*

By applying Proposition 4.2 to the functional $\mathcal{S} : C_T \rightarrow (-\infty, \infty]$ defined in (2.3) with the group

$$G = \left\{ \sum_{k=1}^N k_i \omega_i e^i : k_i \in \mathbb{Z}, i = 1, \dots, N \right\}, \quad (4.1)$$

one obtains easily the following multiplicity result.

Theorem 4.3. *If F satisfies assumption (1.3), then, for each $e \in C([0, T], \mathbb{R}^N)$ such that $\bar{e} = 0$, problem (2.2) has at least $N + 1$ geometrically distinct solutions.*

The proof of Proposition 4.2 given in [1] is quite complicated and technical, but only uses classical Ljusternik-Schnirelmann category. In the following section, we describe a recent approach given in [13], which answers positively the two questions of the end of Section 3, by obtaining the requested multiplicity result through the use of a modified equivalent problem, whose action functional is defined in the classical Sobolev space H_T^1 , and to which Theorem 4.12 of [22] can be directly applied.

5. A MODIFIED LAGRANGIAN APPROACH

Let $e \in L^1(0, T; \mathbb{R}^N)$, and assume that $F(t, \cdot)$ and $\nabla_u F(t, \cdot)$ are continuous for a.e. $t \in [0, T]$, that $F(\cdot, u)$ and $\nabla_u F(\cdot, u)$ are measurable for each $u \in \mathbb{R}^N$, and that there exists some $\alpha \in L^1(0, T)$ such that

$$|F(t, u)| + |\nabla_u F(t, u)| \leq \alpha(t)$$

for a.e. $t \in [0, T]$ and all $u \in \mathbb{R}^N$. Define, for $v \in B(1) \subset \mathbb{R}^N$,

$$\varphi(v) := \frac{v}{\sqrt{1 - |v|^2}},$$

so that

$$\varphi^{-1}(w) = \frac{w}{\sqrt{1 + |w|^2}} \quad \text{for all } w \in \mathbb{R}^N.$$

Let us introduce a modification of φ inspired by recent papers of Coelho *et al* [8, 9] in problems of positive solutions with Dirichlet conditions, but technically different, by setting

$$K := \varphi^{-1}(\bar{B}(\sqrt{n}\|\alpha\|_{L^1})) \subset B(1),$$

fixing $R \in (0, 1)$ in such a way that

$$\frac{R}{\sqrt{1 - R^2}} \geq \sqrt{n}\|\alpha\|_{L^1}, \quad K \subset \bar{B}(R),$$

and defining the homeomorphism $\psi : \mathbb{R}^N \rightarrow \mathbb{R}^N$ by

$$\psi(y) := (1 - \min\{|y|^2, R^2\})^{-1/2}y,$$

in such a way that

$$\psi^{-1}(v) = \max \left\{ (1 - R^2)^{1/2}, (1 + |v|^2)^{-1/2} \right\} v.$$

Lemma 5.1. *For all $y, z \in \mathbb{R}^N$, one has*

$$(\psi(z) - \psi(y)|z - y) \geq |z - y|^2, \quad |\psi(y)| \leq \frac{1}{\sqrt{1 - R^2}}|y|.$$

With this ψ , let us consider the modified problem

$$(\psi(u'))' + \nabla_u F(t, u) = e(t), \quad u(0) = u(T), \quad u'(0) = u'(T) \quad (5.1)$$

The choice of R above allows us to prove the following equivalence result.

Lemma 5.2. *$u \in C^1$ is solution of (2.2) if and only if it is a solution of (5.1).*

If we define now $\Psi : \mathbb{R}^N \rightarrow \mathbb{R}$ by

$$\Psi(y) := 1 - \frac{1 - \min\{|y|^2, R^2\} + 1 - |y|^2}{2\sqrt{1 - \min\{|y|^2, R^2\}}},$$

it is easy to show that Ψ is of class C^1 and that

$$\psi(y) = \nabla \Psi(y), \quad (1/2)|y|^2 \leq \Psi(y) \leq \frac{1}{\sqrt{1 - R^2}}|y|^2$$

for all $y \in \mathbb{R}^N$. Consequently the functional \mathcal{M} given by

$$\mathcal{M}(u) := \int_0^T [\Psi(u') - F(t, u) + (e|u)] dt$$

is well defined and of class C^1 on H_T^1 , and

$$\langle \mathcal{M}'(u), v \rangle = \int_0^T [(\psi(u')|v') - (\nabla_u F(t, u) - e|v)] dt,$$

for all $u, v \in H_T^1$, so that its critical points correspond to the weak, and hence to the Carathéodory solutions of (5.1). On the other hand, the following version of Palais-Smale condition can be proved.

Lemma 5.3. *Each sequence (u_n) in H_T^1 such that $(\mathcal{M}(u_n))$ is bounded, $\mathcal{M}'(u_n) \rightarrow 0$, and (\bar{u}_n) is bounded, contains a convergent subsequence.*

As mentioned above, [22, Theorem 4.12] is just the version of Proposition 4.2 for a C^1 functional, and we keep the notations of Section 4. If $\mathcal{I} \in C^1(X, \mathbb{R})$ is G -invariant, we introduce the following type of Palais-Smale condition, called the $(PS)_G$ -condition : for each sequence (u_n) in X with $(\mathcal{G}(u_n))$ bounded and $\mathcal{G}'(u_n) \rightarrow 0$, $(\pi(u_n))$ contains a convergent subsequence. Theorem 4.12 in [22] goes as follows.

Proposition 5.4. *If the vector space spanned in X by G has finite dimension N , and if $\mathcal{I} \in C^1(X, \mathbb{R})$ is G -invariant, satisfies $(PS)_G$ -condition, and is bounded from below, then \mathcal{I} has at least $N + 1$ critical orbits.*

The proof of Proposition 5.4 given in [22] is based upon Ekeland variational principle and classical Ljusternik-Schnirelmann category. As shown in [22], it provides a proof of Theorem 1.3 for the classical pendulum system. It also implies the corresponding result for (2.2).

Theorem 5.5. *If F satisfies condition (1.3), then, for each $e \in L^1(0, T; \mathbb{R}^N)$, problem (2.2) has at least $N + 1$ geometrically distinct solutions.*

Sketch of the proof. By Lemma 5.2, it suffices to prove that the modified action function \mathcal{M} satisfies the conditions of Proposition 5.4. It is easy to see that

$$\mathcal{M}(u) = \int_0^T [\Psi(u') - F(t, u) + (e(t)|\tilde{u})] dt = \mathcal{M}(u + \omega_i e^i)$$

for all $i = 1, \dots, n$ and $u \in H_T^1$. If we define G by (4.1), using Lemma 5.1 and the Wirtinger and Sobolev inequalities, and denoting the L^p -norm by $\|\cdot\|_p$, one can show that the inequality

$$\mathcal{M}(u) \geq \frac{1}{2}\|\tilde{u}'\|_2^2 - C_2\|\alpha\|_1 - C_1\|h\|_1\|\tilde{u}'\|_2$$

holds, and that \mathcal{M} satisfies the $(PS)_G$ -condition. Then the result follows from Proposition 5.4. \square

REFERENCES

- [1] C. Bereanu, P. Jebelean; *Multiple critical points for a class of periodic lower semicontinuous functionals*, Discrete Continuous Dynam. Syst. A **33** (2013), 47–66.
- [2] C. Bereanu, P. Jebelean, J. Mawhin; *Multiple solutions for Neumann and periodic problems with singular ϕ -Laplacian*, J. Functional Anal. **261** (2011), 3226–3246.
- [3] C. Bereanu, P. Torres; *Existence of at least two periodic solutions for the forced relativistic pendulum*, Proc. Amer. Math. Soc. **140** (2012), 2713–2719.
- [4] H. Brezis, J. Mawhin; *Periodic solutions of the forced relativistic pendulum*, Differential Integral Equations **23** (2010), 801–810.
- [5] H. Brezis, J. Mawhin; *Periodic solutions of Lagrangian systems of relativistic oscillators*, Comm. Appl. Anal. **15** (2011), 235–250.
- [6] A. Castro; *Periodic solutions of the forced pendulum equations*, in Differential Equations (Oklahoma, 1979), Ahmad-Lazer (ed.), Academic Press, New York, 1980, 149–160.
- [7] K. C. Chang; *On the periodic nonlinearity and the multiplicity of solutions*, Nonlinear Anal. **13** (1989), 527–537.
- [8] I. Coelho, C. Corsato, F. Obersbel, P. Omari; *Positive solutions of the Dirichlet problem for the one-dimensional Minkowski-curvature equation*, Advanced Nonlinear Studies **12** (2012), 621–638.
- [9] I. Coelho, C. Corsato, S. Rivetti; *Positive radial solutions of the Dirichlet problem for the Minkowski-curvature equation in a ball*, Topological Methods Nonlinear Anal. **44** (2014), 23–40.
- [10] E. N. Dancer; *On the use of asymptotics in nonlinear boundary value problems*, Ann. Mat. Pura Appl. **131** (1982), 167–187.
- [11] A. Fonda, R. Toader; *Periodic solutions of pendulum-like Hamiltonian systems in the plane*, Advanced Nonlinear Studies **12** (2012), 395–408.
- [12] G. Hamel; *Ueber erzwungene Schwingungen bei endlichen Amplituden*, Math. Ann. **86** (1922), 1–13.
- [13] P. Jebelean, J. Mawhin, C. Serban; *Multiple periodic solutions for perturbed relativistic pendulum systems*, Proc. Amer. Math. Soc., 143 (2015), 3029–3039.
- [14] R. Manásevich, J.R. Ward Jr; *On a result of Brezis and Mawhin*, Proc. Amer. Math. Soc. **140** (2012), 531–539.
- [15] S. Marò; *Periodic solutions of a forced relativistic pendulum via twist dynamics*, Topological Methods Nonlinear Anal. **42** (2013), 51–75.
- [16] J. Mawhin; *The forced pendulum : a paradigm for nonlinear analysis and dynamical systems*, Expositiones Math. **6** (1988), 271–287.
- [17] J. Mawhin; *Forced second order conservative systems with periodic nonlinearities*, Ann. Inst. Henri-Poincaré Anal. non linéaire **6** (1989), suppl., 415–434.
- [18] J. Mawhin; *Global results for the forced pendulum equation*, in Handbook of Differential Equations, Ordinary Differential Equations, Vol. 1, A. Cañada, P. Drábek, A. Fonda (ed.), Elsevier, Amsterdam, 2004, 533–589.
- [19] J. Mawhin; *Multiplicity of solutions of variational systems involving ϕ Laplacians with singular ϕ and periodic nonlinearities*, Discrete Continuous Dynam. Syst. A **32** (2012), 4015–4026.
- [20] J. Mawhin, M. Willem; *Multiple solutions of the periodic boundary value problem for some forced pendulum-type equations*, J. Differential Equations **52** (1984), 264–287.
- [21] J. Mawhin, M. Willem; *Variational methods and boundary value problems for vector second order differential equations and applications to the pendulum equation*, in Nonlinear Analysis and Optimization, Lect. Notes Math. **1107**, Springer, Berlin, 1984, 181–192.

- [22] J. Mawhin, M. Willem; *Critical Point Theory and Hamiltonian Systems*, Springer, New York, 1989.
- [23] P. H. Rabinowitz; *Minimax Methods in Critical Point Theory with Applications to Differential Equations*, CMBS Reg. Conf. Math. **65**, Amer. Math. Soc., Providence RI, 1986.
- [24] P. H. Rabinowitz; *On a class of functionals invariant under a Z^n action*, Trans. Amer. Math. Soc. **310** (1988), 303–311.
- [25] A. Szulkin; *Minimax principles for lower semicontinuous functions and applications to non-linear boundary value problems*, Ann. Inst. Henri-Poincaré Anal. non lin. **3** (1986), 77–109.
- [26] A. Szulkin; *A relative category and applications to critical point theory for strongly indefinite functionals*, Nonlinear Anal. **15** (1990), 725–739.
- [27] M. Willem; *Oscillations forcées de l'équation du pendule*, Publ. IRMA Lille **3** (1981), v1–v3.

JEAN MAWHIN

INSTITUT DE RECHERCHE EN MATHÉMATIQUE ET PHYSIQUE, UNIVERSITÉ CATHOLIQUE DE LOUVAIN,
1348 LOUVAIN-LA-NEUVE, BELGIUM

E-mail address: jean.mawhin@uclouvain.be